


# thermodynamics of effective stress in partially saturated soils

By: Mario Liu, Yimin Jiang and Itai Einav



Universität Tübingen,  
Germany



Central South University,  
China

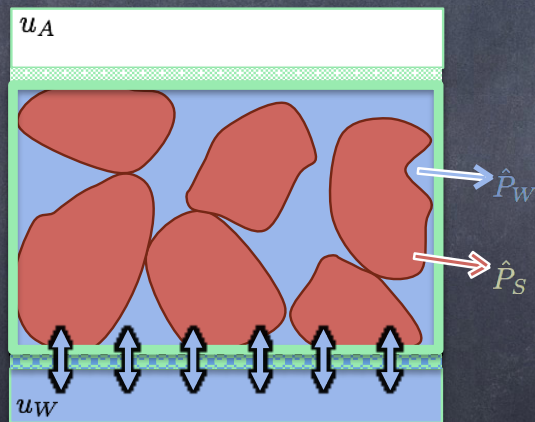


University of Sydney,  
Australia

# Terzaghi's effective stress

$$\sigma_{ij}^{eff} = \sigma_{ij} - P_T \delta_{ij}$$

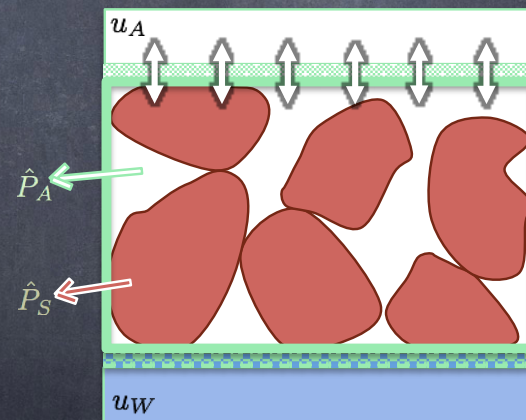
$S_r=1$ , fully wet



$$u_W = \hat{P}_W = \hat{P}_S = u_A$$
$$s = u_A - u_W = 0$$

$$P_T = u_W$$

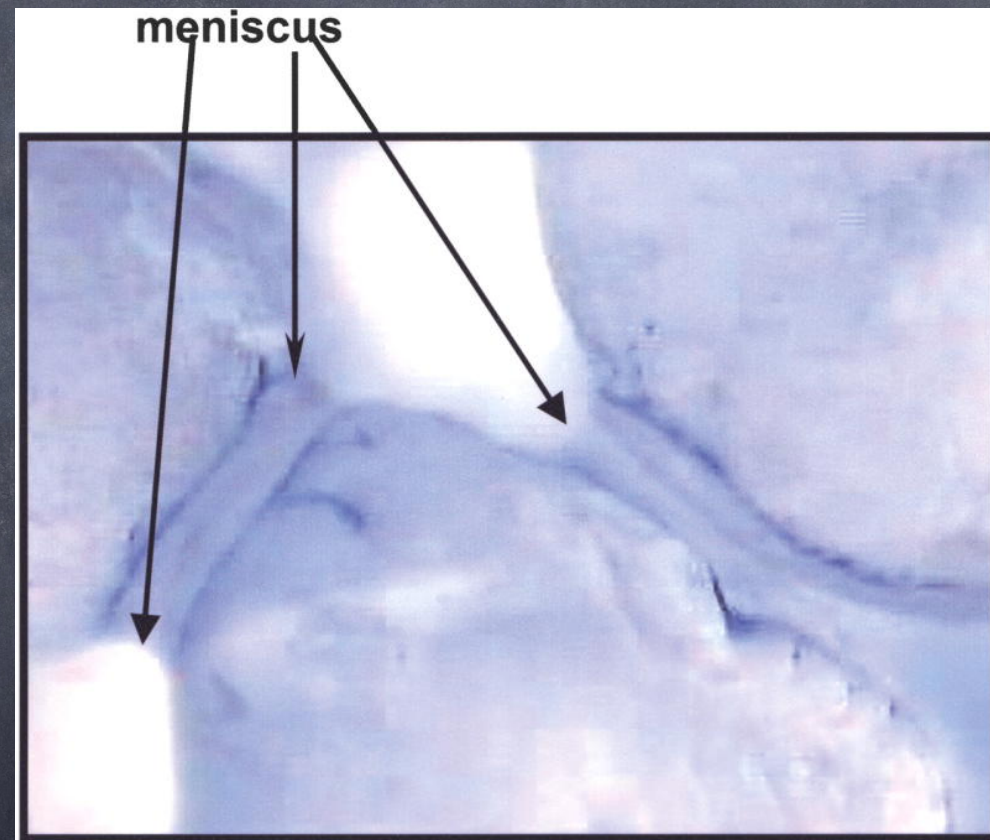
$S_r=0$ , fully dry



$$u_A = \hat{P}_A = \hat{P}_S = u_W$$
$$s = u_A - u_W = 0$$

$$P_T = u_A$$

# partially saturated media



Solid-Air-Water interfaces

# thermodynamic treatments

- previous treatments try to follow microscopic surface tensions of complex shaped interfaces
- current treatment follows macroscopic suction from soil-water retention curves

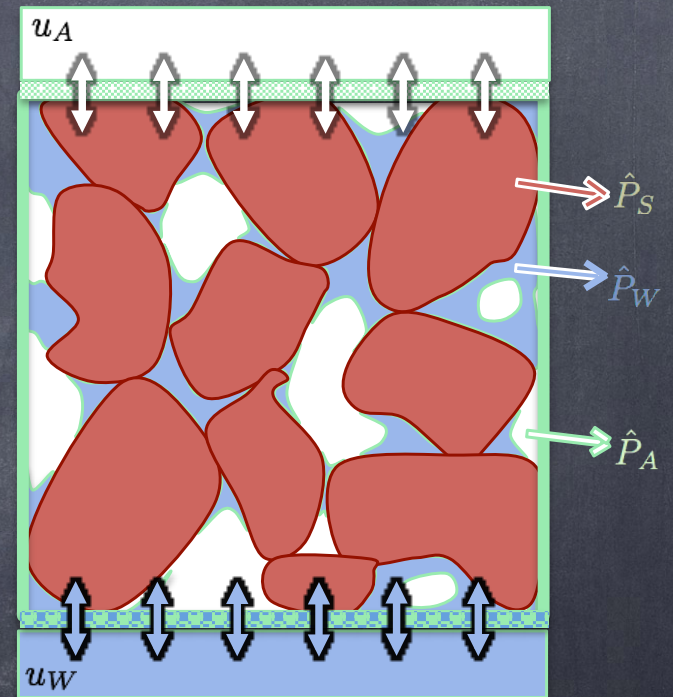
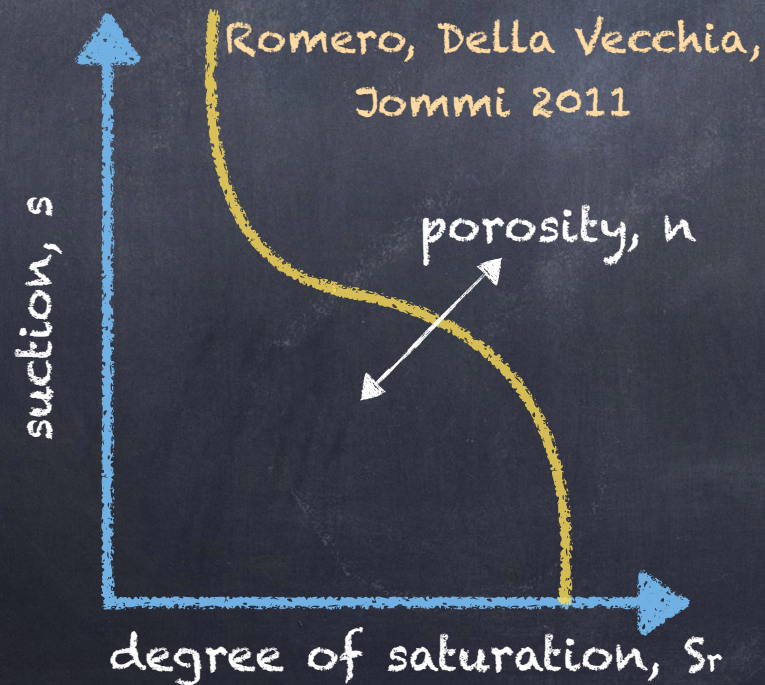
Note, soil-water retention curves encode the effects of surface tensions of complex interfaces

# partially saturated media ( $1 > S_r > 0$ )

$$\phi_\beta \equiv \frac{V_\beta}{V}, \quad \sum_\beta \phi_\beta = 1, \quad \beta = (A, W, S),$$

$$n \equiv \phi_W + \phi_A, \quad S_r \equiv \frac{\phi_W}{n}.$$

soil-water retention curves



$$s = u_A - u_W$$

$$\stackrel{?}{=} \hat{P}_A - \hat{P}_W$$

# Bishop's effective stress

$$\sigma_{ij}^{eff} = \sigma_{ij} - P_T \delta_{ij}$$

partially saturated soils ( $1 > S_r > 0$ )

$$P_T = u_A - \chi (u_A - u_W) \stackrel{=s}{=}$$

$$\chi \equiv \chi(S_r) \dots \text{or} \dots \chi \equiv \chi(s) \dots \text{or} \dots ???$$

but soil-water retention shows  $s = s(S_r, n)$

# methodology

- develop a hydrodynamic treatment without needing to assume the shapes of interfaces
- uncover the meaning of pressure  $P_T$
- distinguish intrinsic and partial quantities
- reveal the structure of effective stress

# partial quantities

$\rho_\beta \equiv M_\beta/V$ ,  $\beta = (A, W, S)$  = partial densities

$\rho \equiv \sum \rho_\beta$  = total density

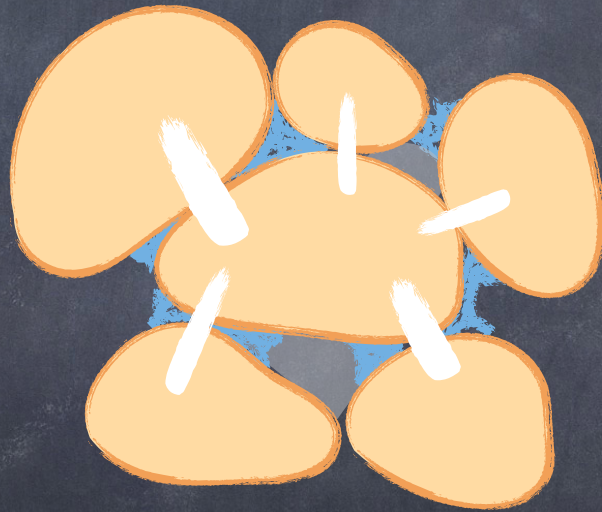
$f \equiv f(\rho_\beta) = f(\rho_A, \rho_W, \rho_S)$  = free energy

$\mu \equiv \partial f / \partial \rho$  = chemical potential (Gibbs, 1878)



# general hydrodynamic result

$$\sigma_{ij}^e = \sigma_{ij} - P_T \delta_{ij}$$



$$P_T = -\frac{\partial(f/\varrho)}{\partial(1/\varrho)} = \boxed{P_T \equiv \varrho\mu - f} \text{ = thermodynamic pressure}$$

# intrinsic quantities

$$\hat{f}_\beta \equiv \hat{f}_\beta(\hat{\varrho}_\beta), \quad = \text{intrinsic free}$$

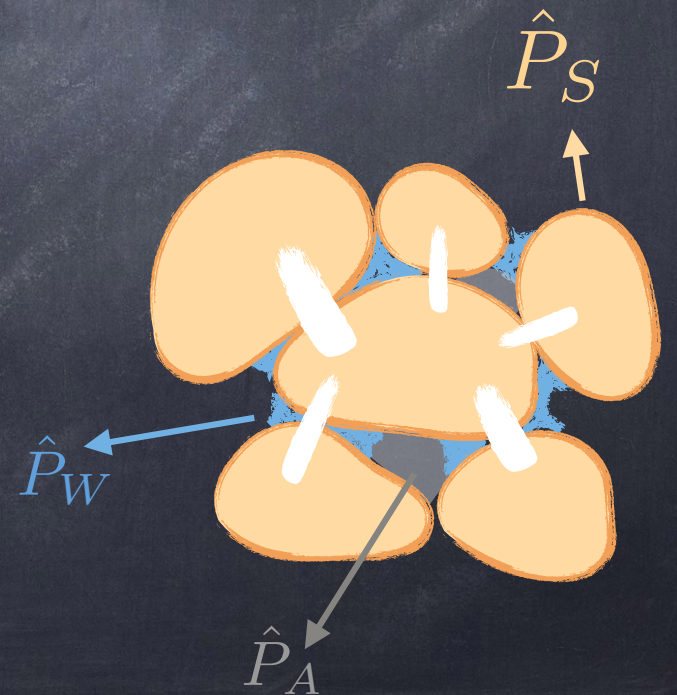
$$\hat{\mu}_\beta \equiv \partial \hat{f}_\beta / \partial \hat{\varrho}_\beta, \quad = \text{intrinsic chemical}$$

$$\hat{P}_\beta \equiv \hat{\varrho}_\beta \hat{\mu}_\beta - \hat{f}_\beta. \quad = \text{intrinsic pressure}$$

$$f = \sum \phi_\beta \hat{f}_\beta(\hat{\varrho}_\beta) \quad = \text{total free}$$

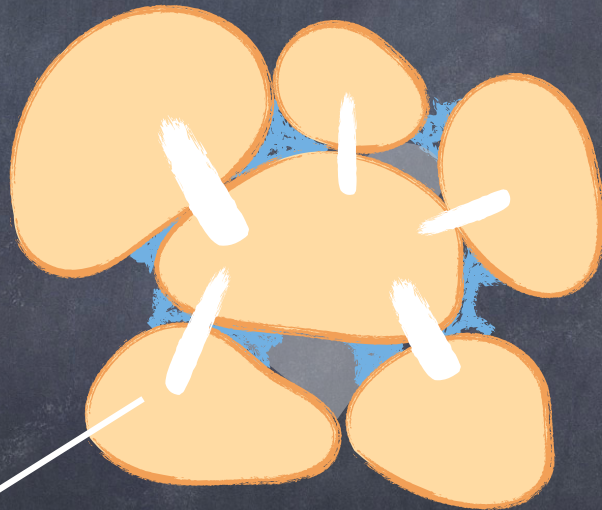
(total) thermodynamic pressure

$$P_T = \sum_\beta \hat{P}_\beta \left[ 1 - \rho \frac{\partial}{\partial \rho} \right] \phi_\beta$$



# effective stress and elastic forces

$$\sigma_{ij}^{eff} = \sigma_{ij} - P_T \delta_{ij}$$



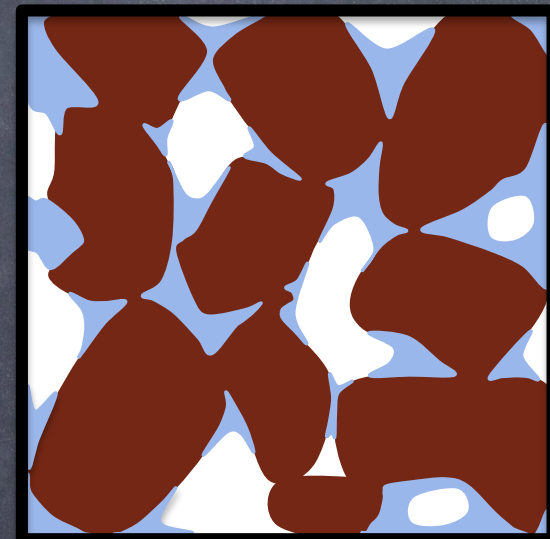
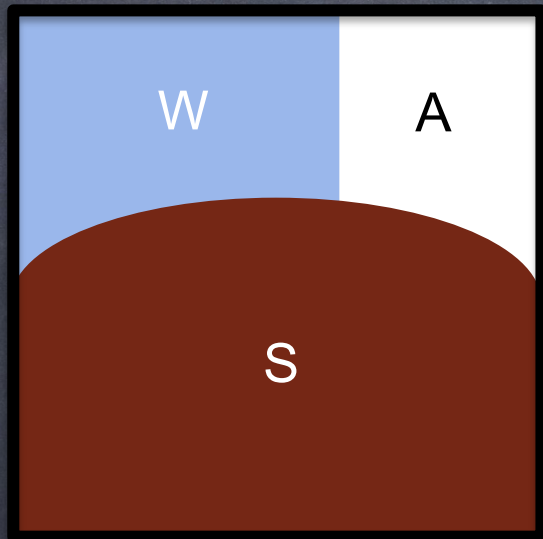
$$\sigma_{ij}^{eff} = \sigma_{ij}^e$$

(elastic)  
contact  
forces

$$\rightarrow P^{eff} = \frac{1}{3} \sigma_{ii}^{eff}$$

use to  
integrate  
constitutive  
relations

# 'suctionless limit' and 'common pressure'



$$P_T = P_0 = \hat{P}_\beta$$

$$\Delta \hat{P}_\beta = \hat{P}_\beta - P_0 = K_\beta \frac{\Delta \hat{Q}_\beta}{\hat{Q}_\beta}$$

# minimisation of free energy

$$\delta \int \left( f - L_1 \sum_{\beta} \phi_{\beta} - L_2 [\hat{P}_A - \hat{P}_W] \right) d^3 r = 0$$

mass constraint

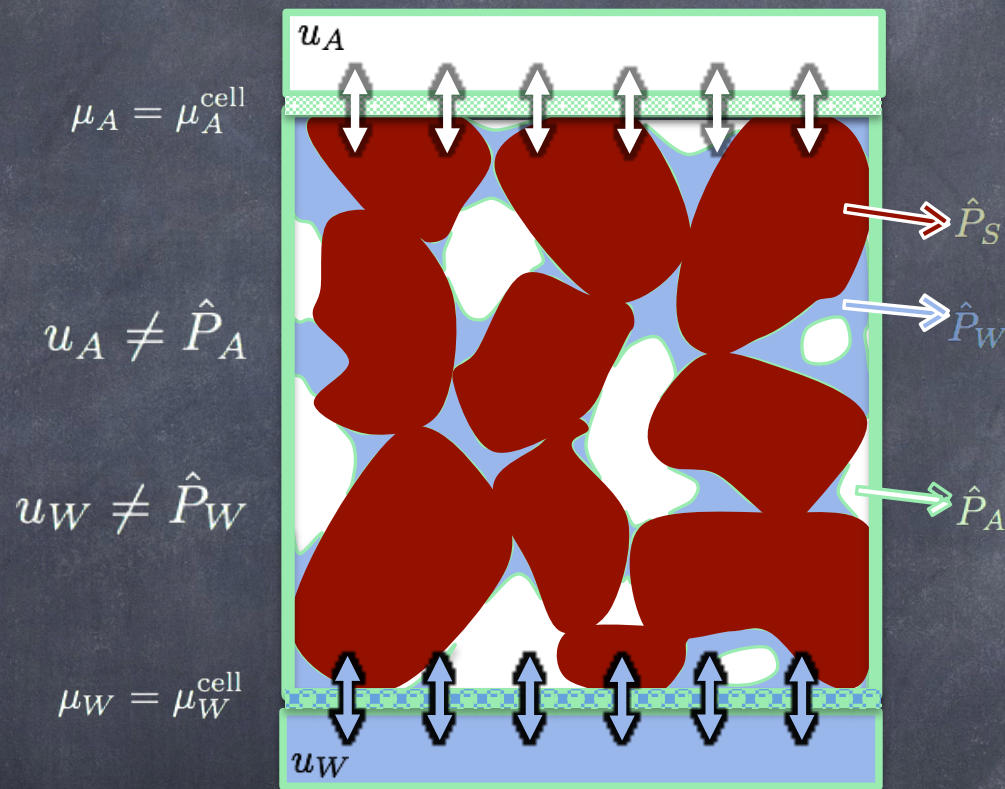
$= \hat{s}$   
suction constraint

$$\hat{s} = L_2 \left( \frac{K_A}{\phi_A} + \frac{K_W}{\phi_W} \right) \quad \text{where} \quad K_{\beta} = \hat{q}_{\beta} \frac{\partial \hat{P}_{\beta}}{\partial \hat{q}_{\beta}}$$

...and by accounting for realistic compressibilities:

$$\begin{aligned} \hat{s} &= \hat{P}_A - \hat{P}_W \\ &= \Delta \hat{P}_A - \Delta \hat{P}_W \end{aligned}$$

# equilibrium and suctions



$$\mu_\beta = \mu_\beta^{\text{cell}}$$

$$s = u_A - u_W \text{ = 'measured suction'}$$

$$\hat{s} = \hat{P}_A - \hat{P}_W \text{ = 'intrinsic suction'}$$

$$\hat{s} \neq s$$

second order approximation  
around suctionless limit

$$\Delta u_\beta = \frac{\phi_W}{2K_W} \hat{\rho}_\beta \psi_\beta, \quad \psi = \hat{s}^2, \quad \psi_\beta = \frac{\partial \psi}{\partial \rho_\beta}$$

therefore...

# effective stress

general structure

$$\hat{s}^2 = \psi(\rho_A, \rho_W, \rho_S)$$

$$P_T = u_A - \chi(u_A - u_w)$$

$$\chi = \chi(\rho_A, \rho_W, \rho_S) = \phi_w + \phi_s \left[ \frac{\hat{\rho}_A \psi_A - \hat{\rho}_S \psi_S}{\hat{\rho}_A \psi_A - \hat{\rho}_W \psi_W} \right]$$

$$s = s(\rho_A, \rho_W, \rho_S) = \frac{\phi_w}{2K_w} (\hat{\rho}_A \psi_A - \hat{\rho}_W \psi_W)$$



# effective stress

SWRC dependent on saturation only

$$\hat{s}^2 = \psi(S_r)$$

$$P_T = u_A - \chi(u_A - u_w)$$

$$\chi = \chi(S_r) = S_r$$

$$s = s(S_r) = -\frac{S_r \psi_{sr}}{2K_w}$$

# effective stress

SWRC dependent on saturation and porosity

$$\hat{s}^2 = \psi(S_r, n)$$

$$P_T = u_A - \chi(u_A - u_w)$$

$$\chi = \chi(S_r, n) = S_r - n(1 - n) \frac{\psi_n}{\psi_{sr}}$$

$$s = s(S_r, n) = -\frac{S_r \psi_{sr}}{2K_W}$$

# effective stress

example

$$\hat{s}^2 = \psi = 2AK_W \phi_S^\beta \left[ \frac{1 - \alpha + \alpha S_r - S_r^\alpha}{\alpha(1 - \alpha)S_r^\alpha} \right]$$

$$P_T = u_A - \chi(u_A - u_w)$$

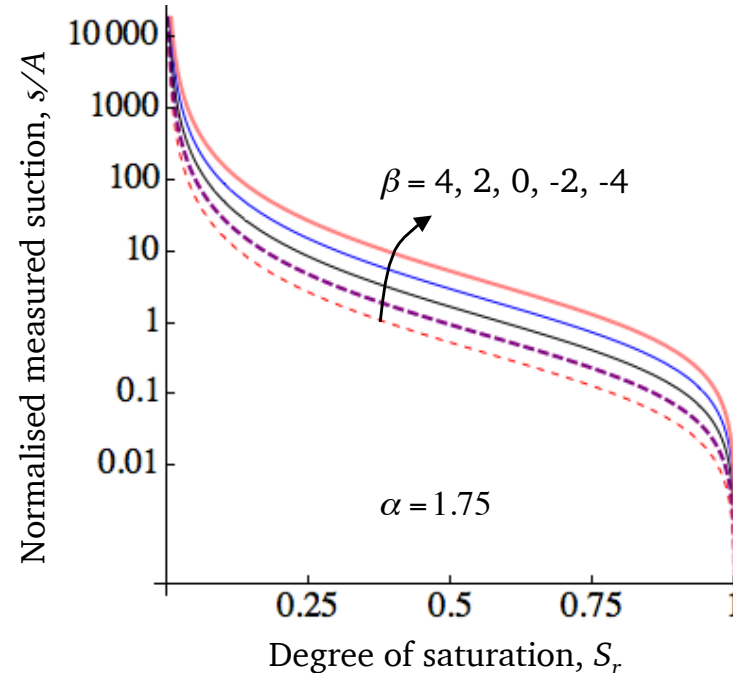
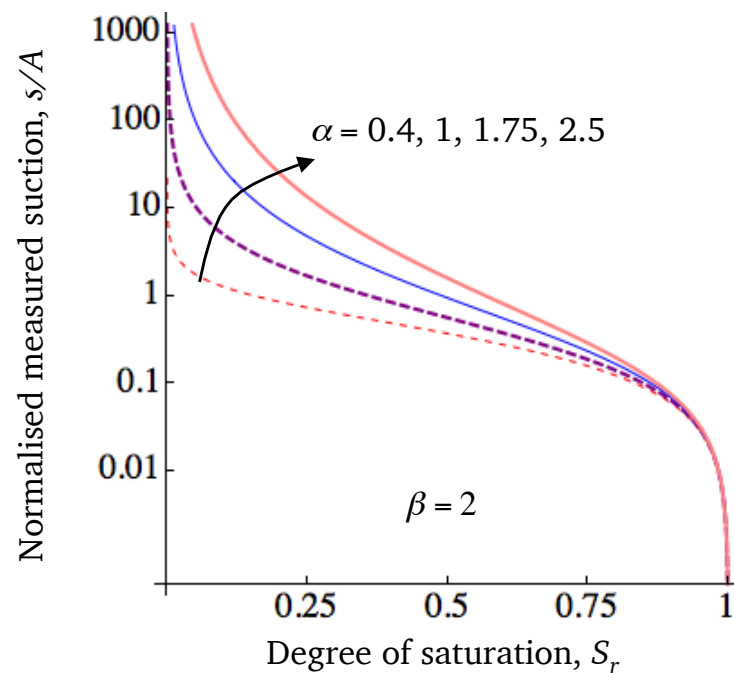
$$\chi = S_r - \beta S_r(1 - n) \left[ \frac{1 - \alpha + \alpha S_r - S_r^\alpha}{\alpha(1 - \alpha)(1 - S_r)} \right]$$

$$s = A(1 - n)^\beta \left[ \frac{1 - S_r}{S_r^\alpha} \right]$$

# effective stress

example

$$s = A(1 - n)^{\beta} \left[ \frac{1 - S_r}{S_r^{\alpha}} \right]$$

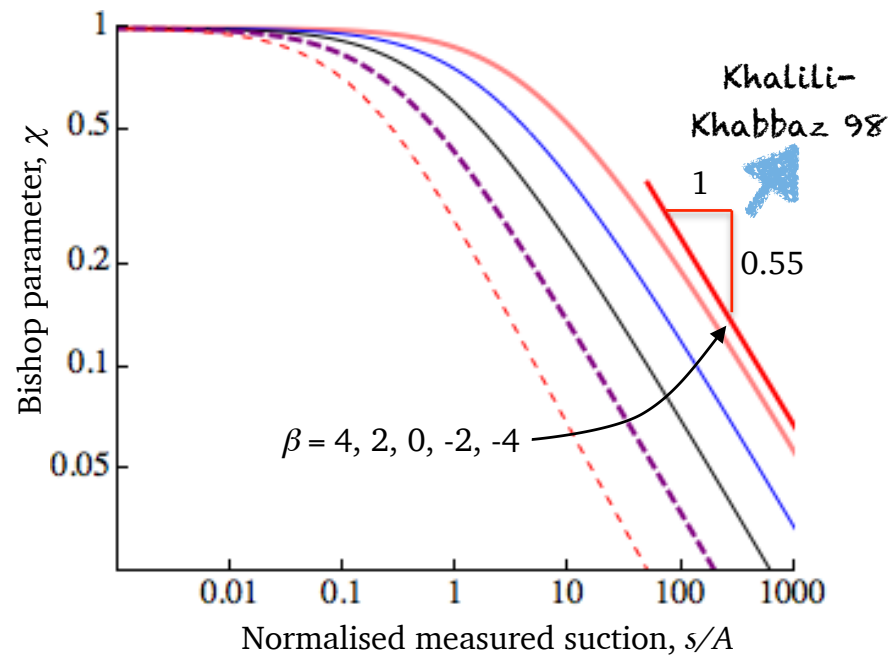
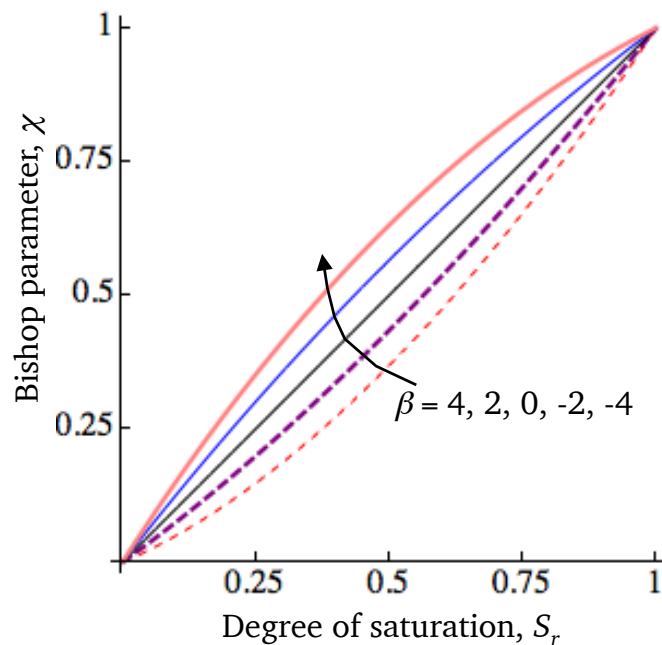


# effective stress

example

$$s = A(1 - n)^\beta \left[ \frac{1 - S_r}{S_r^\alpha} \right]$$

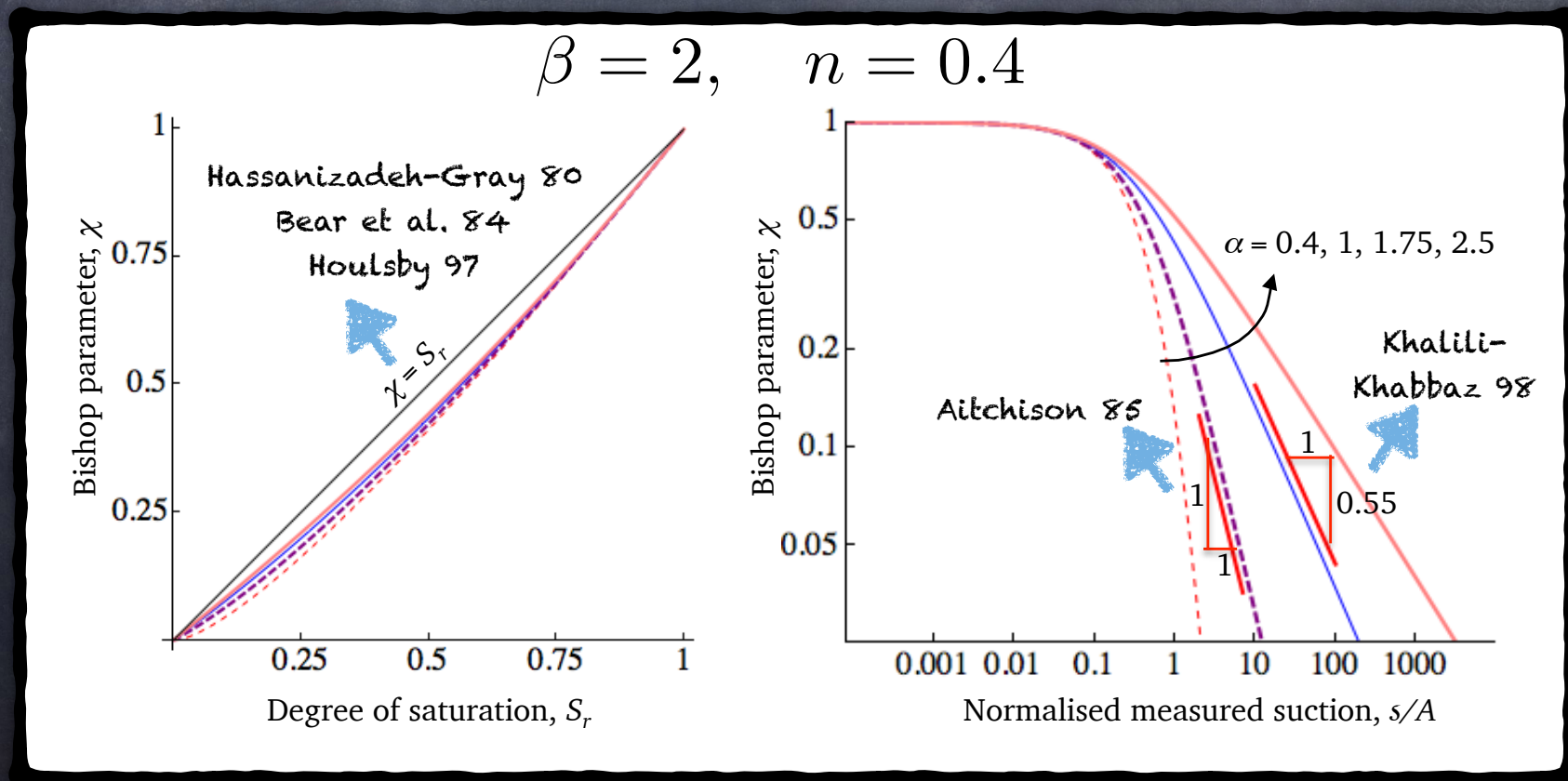
$$\alpha = 1.75, \quad n = 0.4$$



# effective stress

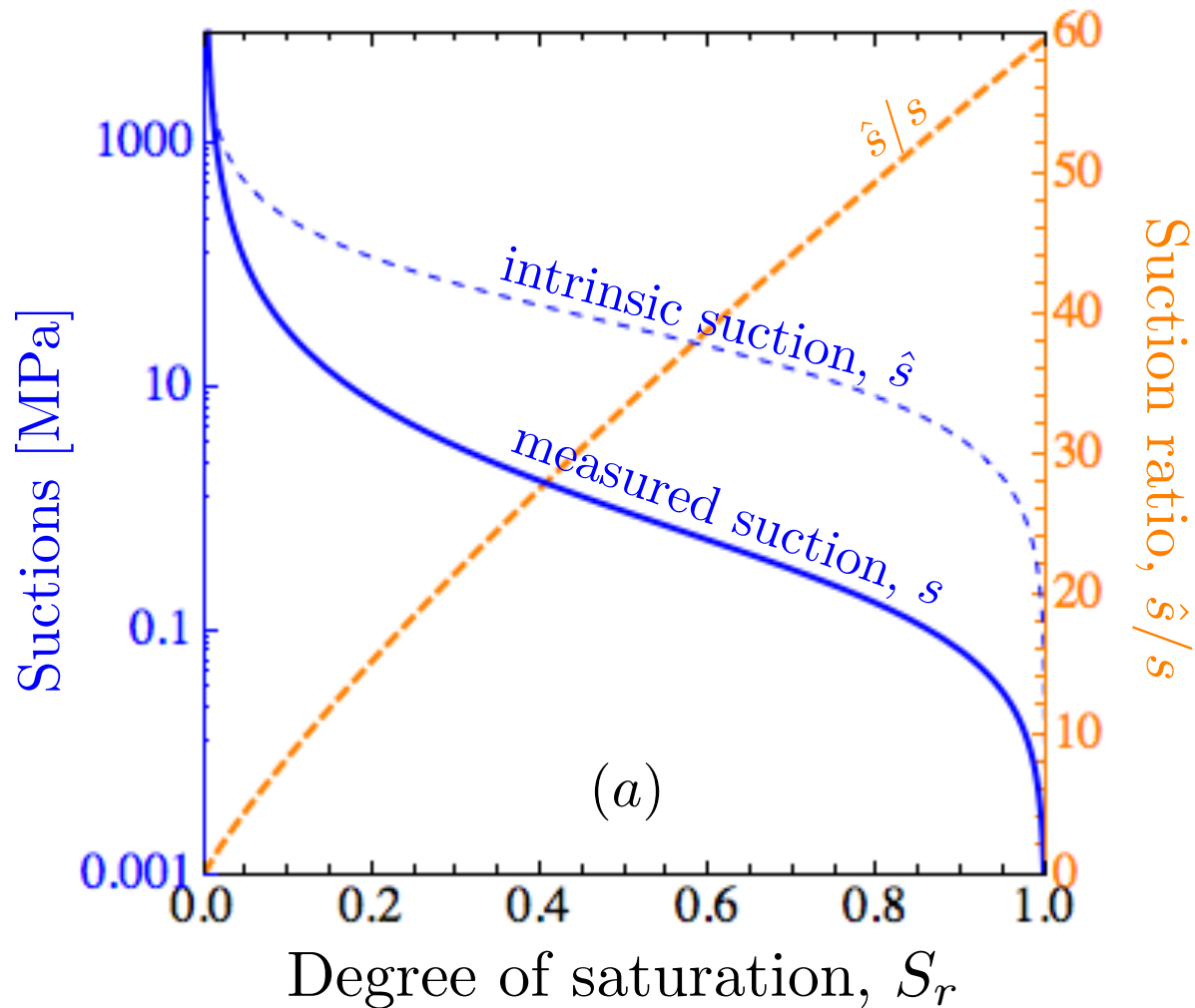
example

$$s = A(1 - n)^\beta \left[ \frac{1 - S_r}{S_r^\alpha} \right]$$



# effective stress

example



# conclusions

- hydrodynamics reveal the structure of effective stress in partially saturated soils
- its structure agrees with Bishop, but the  $\chi$  parameter generally depends on 3 densities in a way strictly connected to soil-water retention
- realistic soil-water retention curves that are a function of both  $n$  and  $S_r$ , explain the versatility of observed experimental  $\chi$  values.



# what's not included (left for future work)

- cohesion effects in very low saturations
- mass exchanges between phases