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Multi-physics Multi-scale Modelling in Geomechanics

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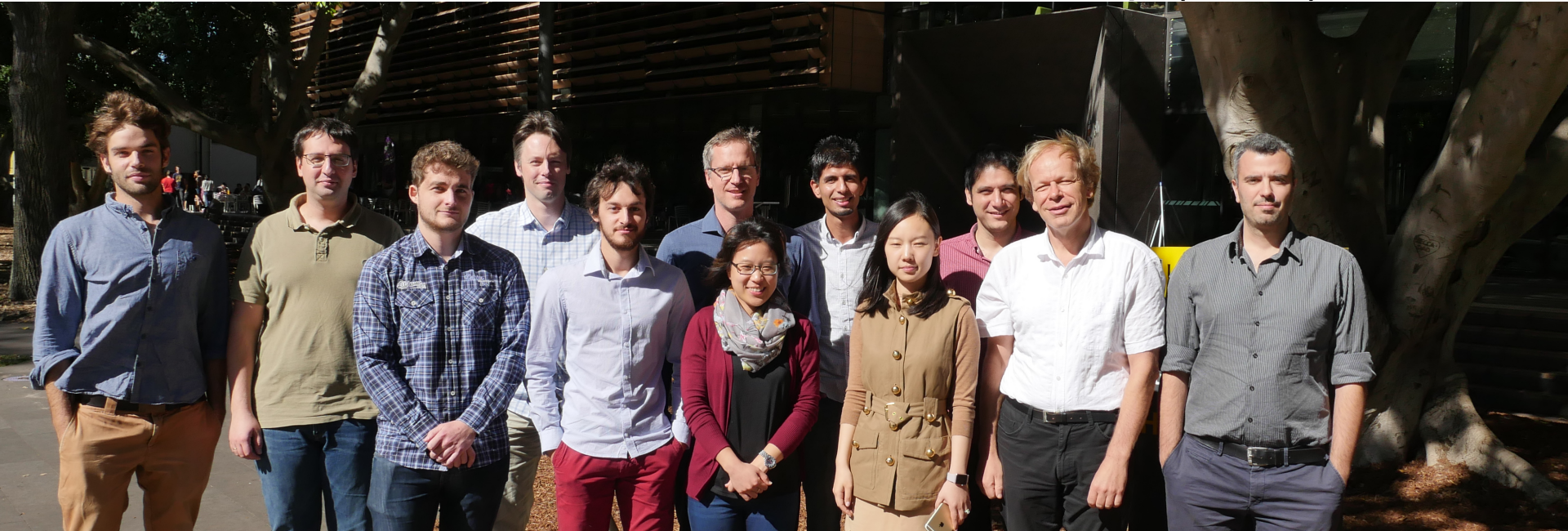
UNSW Petroleum Engineering

CSIRO Energy and Mineral Resources

Never Stand Still

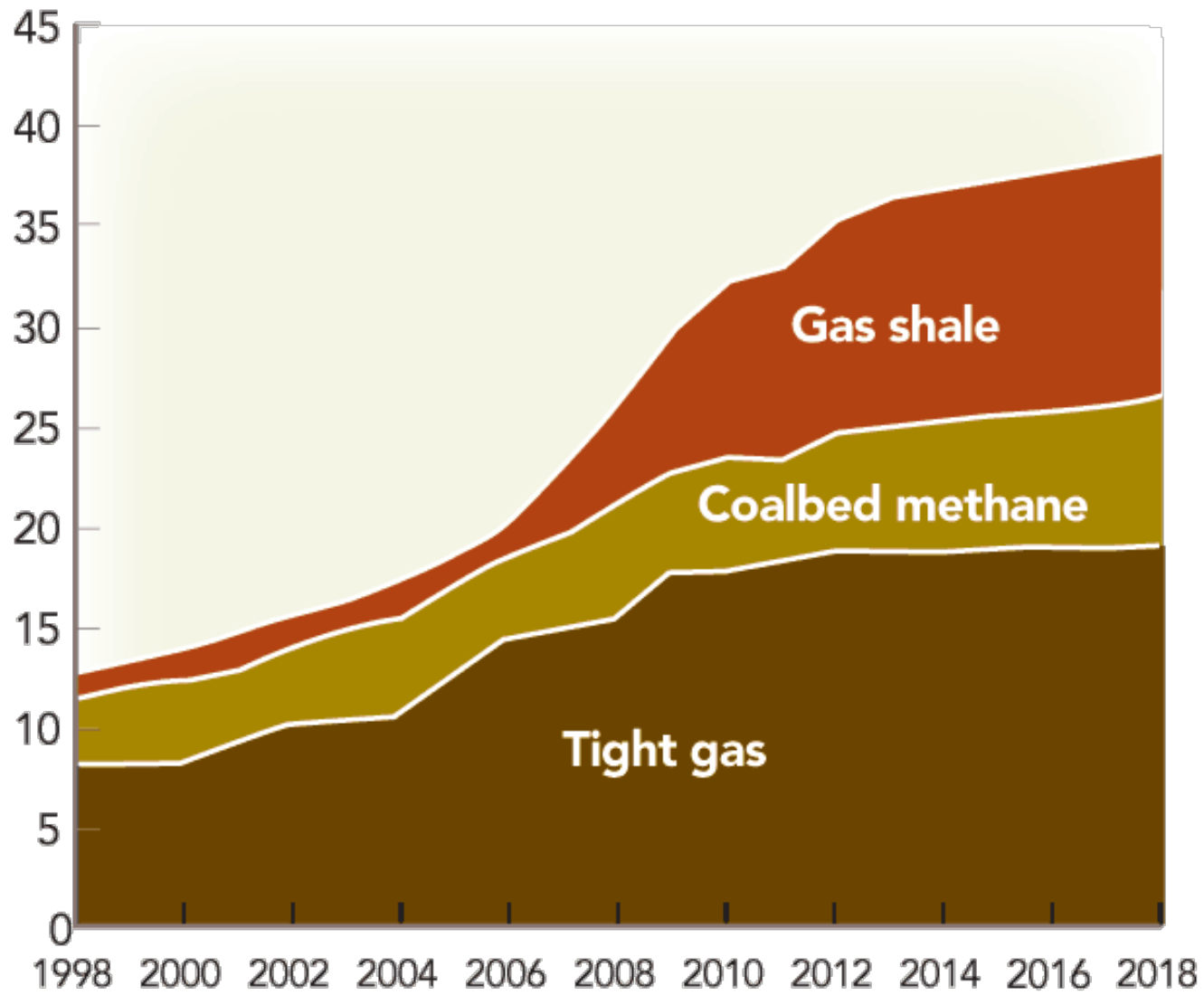
School of Petroleum Engineering

Unconventional Geomechanics Group (UGG)



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First Question: Why Multi-physics?



Source: US DOE 2008

DIAGENETIC SETTINGS

BASED ON MINERALOGY, GEOCHEMISTRY, PETROLEUM, AND HYDROGEOLOGY

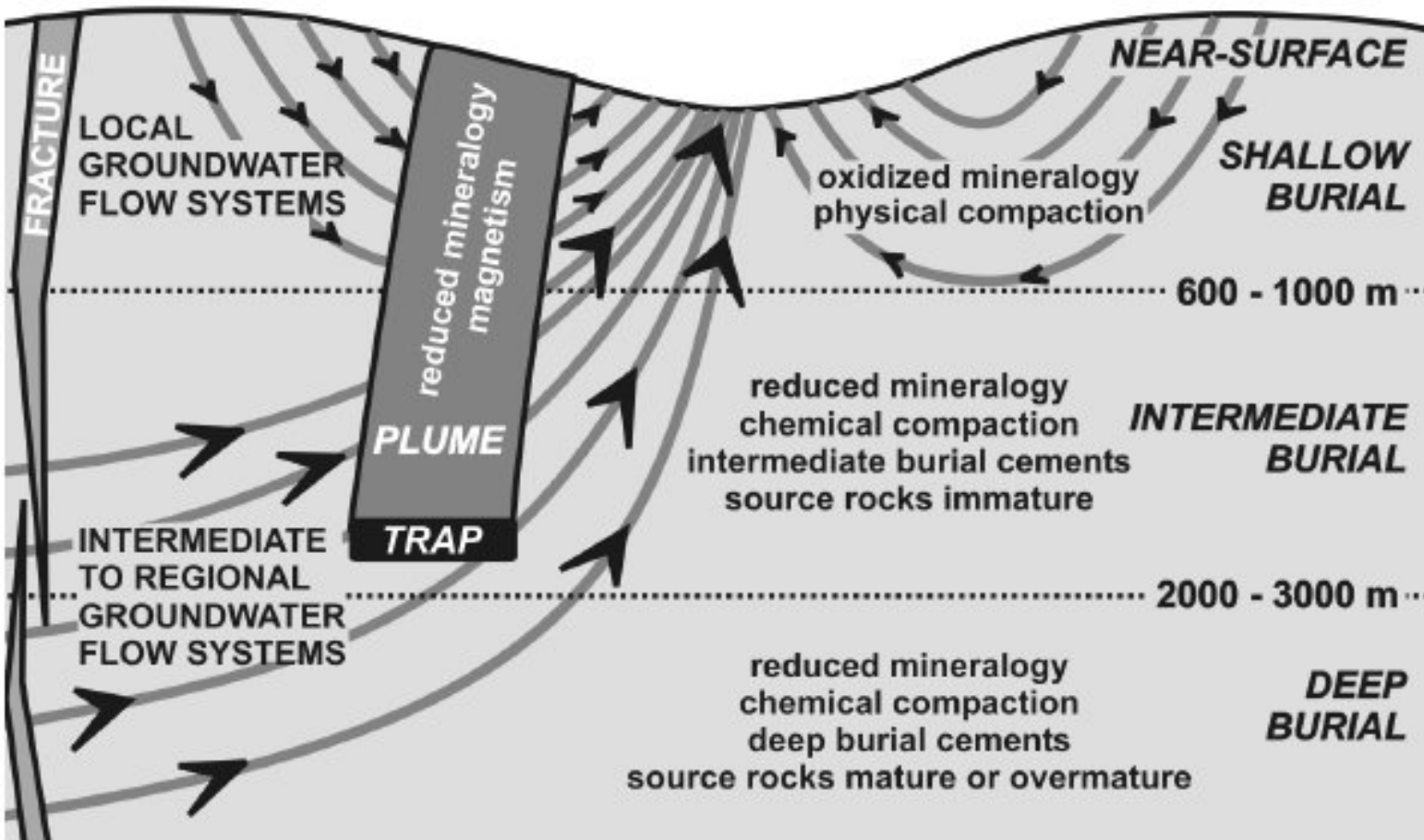
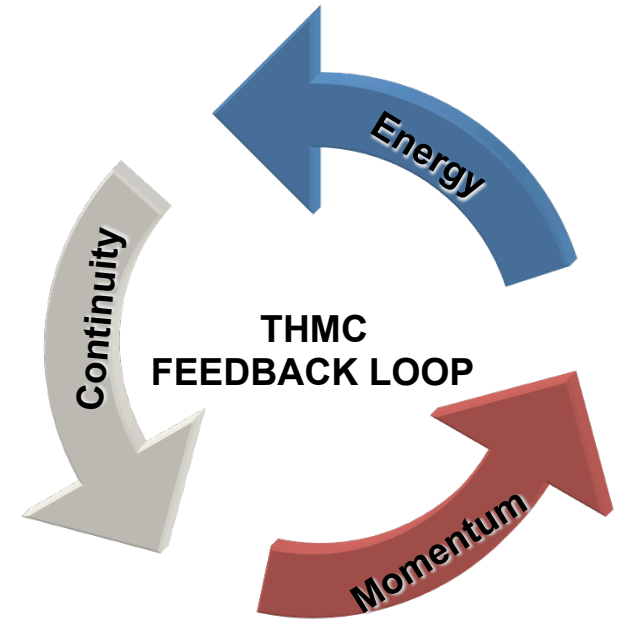
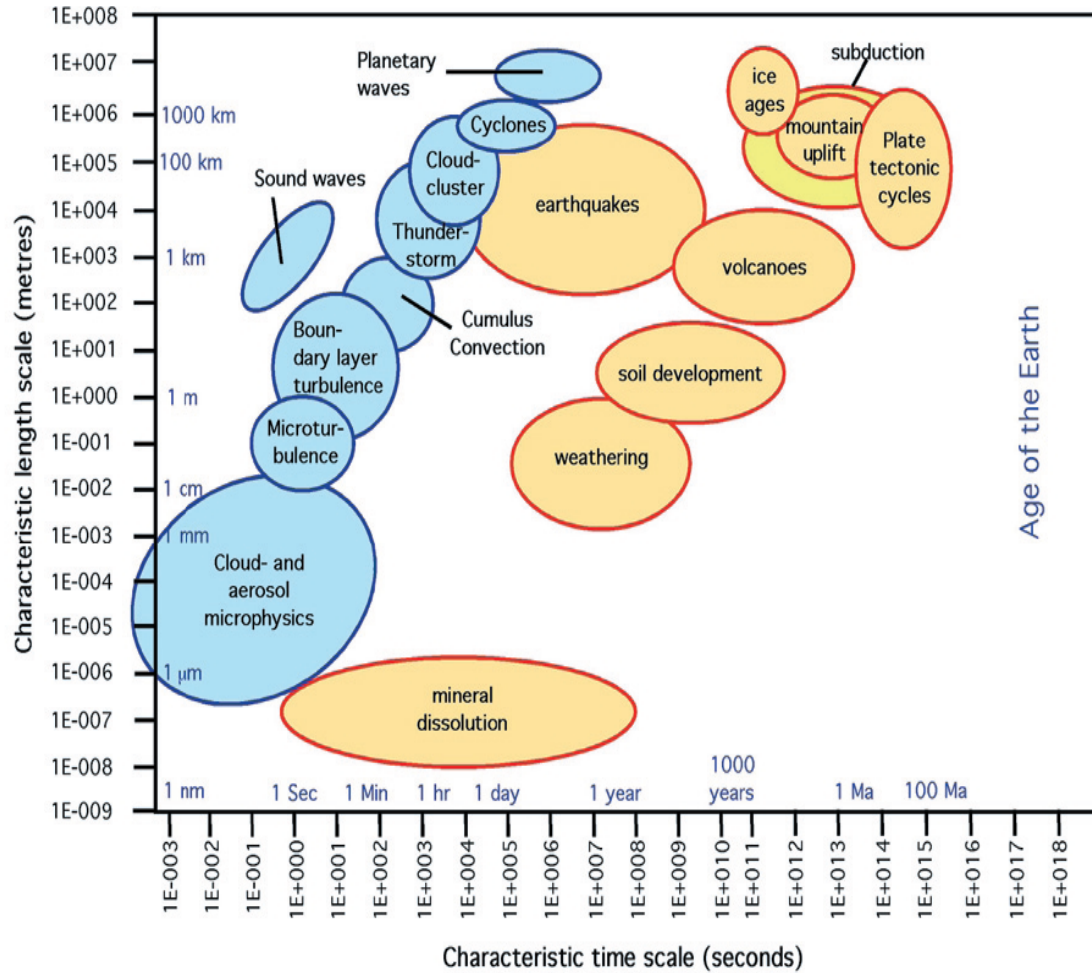
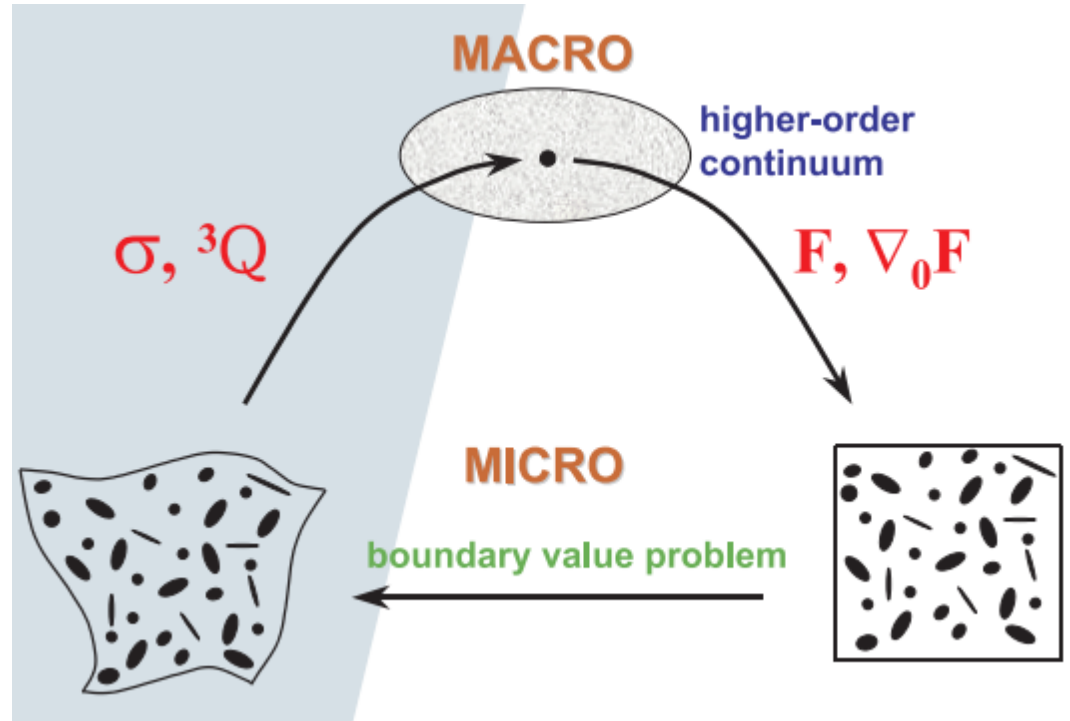


Figure 2. Classification of diagenetic settings on the basis of mineralogy, petroleum, hydrogeochemistry, and hydrogeology. For illustrative simplicity, the geologic section is assumed to be isotropic and homogeneous, with idealized groundwater flow lines. The hydrocarbon-contaminated plume is slightly deflected by the local and regional groundwater flow systems. The depth limits separating the burial diagenetic settings are approximate and based on geologic phenomena that are easily recognizable. Near-surface settings may be meteoric, brackish, marine, or hypersaline. Adapted from Machel (1999).

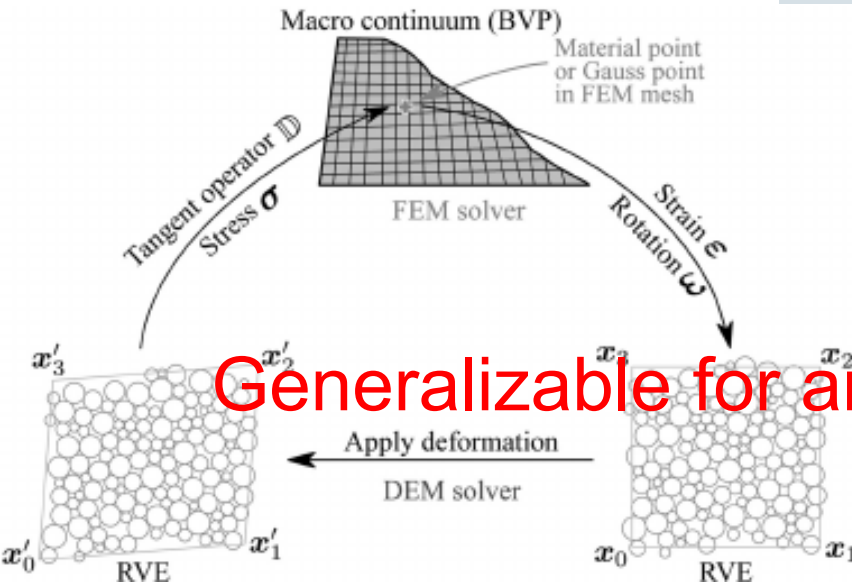
Second Question: Why Multi-scale?



Current Paradigm: computational homogenisation/upscaling



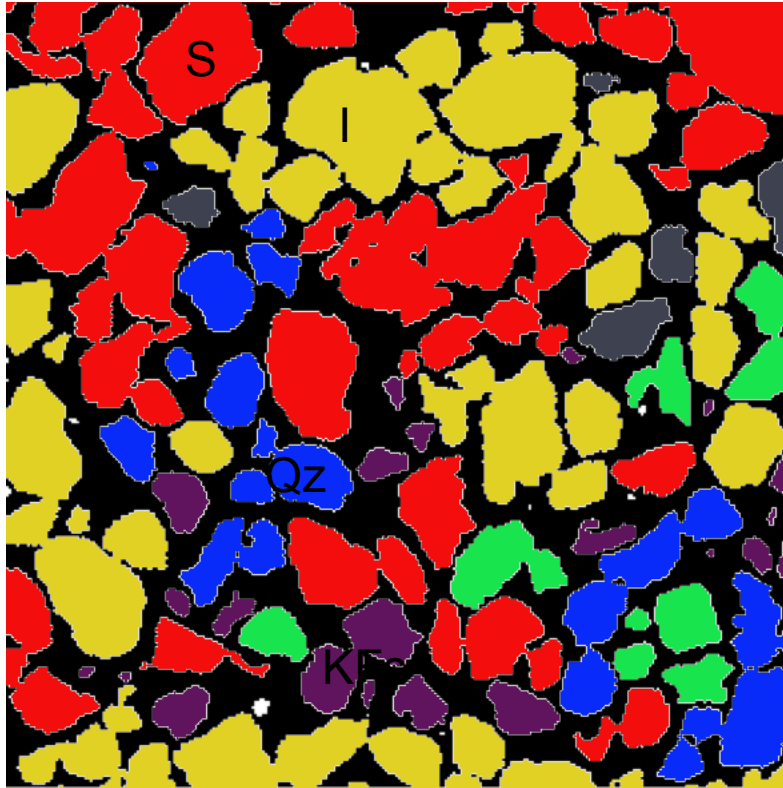
Zhao and Guo, 2011



Dulikravich, G. S., & Tanaka, M. (2000). Inverse Problems in Engineering Mechanics II. Elsevier.

Generalizable for any thermodynamic flux/force

How does it work?



Stack of CT-scan images
Chemical composition

Chemical system reduction (chemical homogenisation)

$$\frac{d[S]}{dt} = -[S]R_1$$

$$\frac{d[X]}{dt} = [S]R_1$$

$$S_i \left(\frac{d[I]}{dt} = f[I][L]R_3 \right.$$

$$K_i \left(\frac{d[K]}{dt} = [F]R_2 - [K][A][X]^f R_3 \right.$$

$$K_i \left(\frac{d[A]}{dt} = [F]R_2 - [K][A][X]^f R_3 \right.$$

$$S_i \left(\frac{d[L]}{dt} = [F]R_2 + [K][A][X]^f R_3 + [L](R_4^- - R_4^+) \right.$$

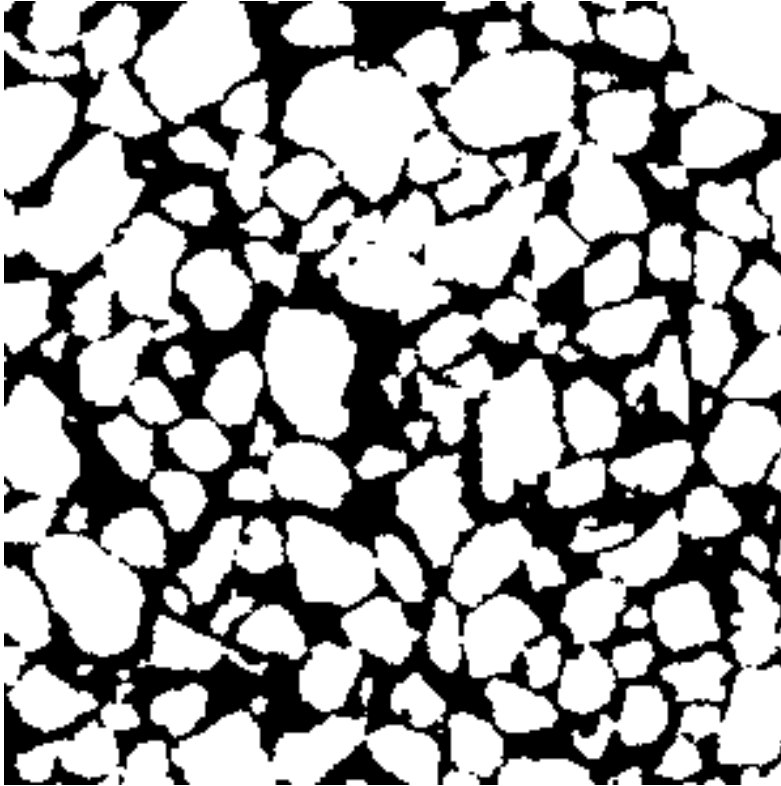
$$\frac{d[Q]}{dt} = [L](R_4^+ - R_4^-)$$

$$\frac{d[F]}{dt} = -[F]R_2$$

$$\frac{d[W]}{dt} = n[S]R_2$$

d)

How does it work?

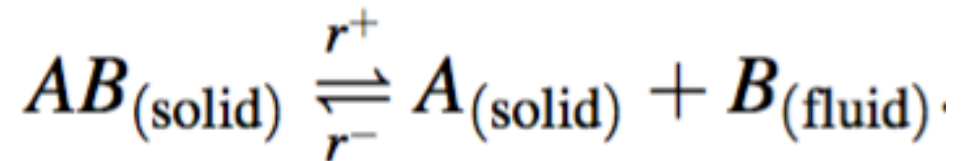


Stack of CT-scan images
Chemical composition

Chemical system reduction (chemical homogenisation)

Invariant manifold selection. E.g. steady state approximation for the intermediate species:

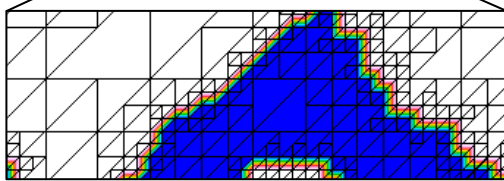
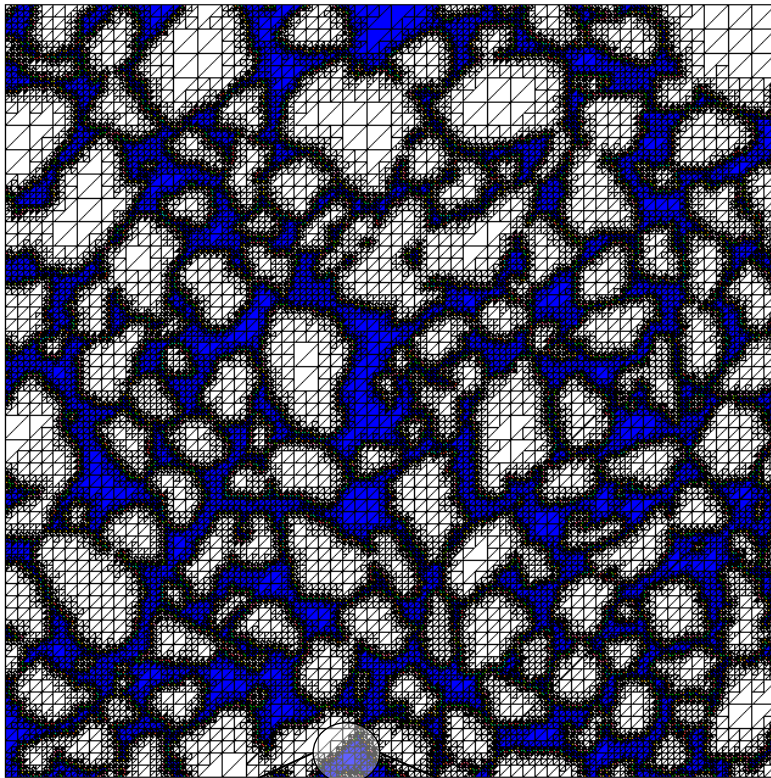
$$\frac{d[X]}{dt} = \frac{d[A]}{dt} = \frac{d[L]}{dt} = 0$$



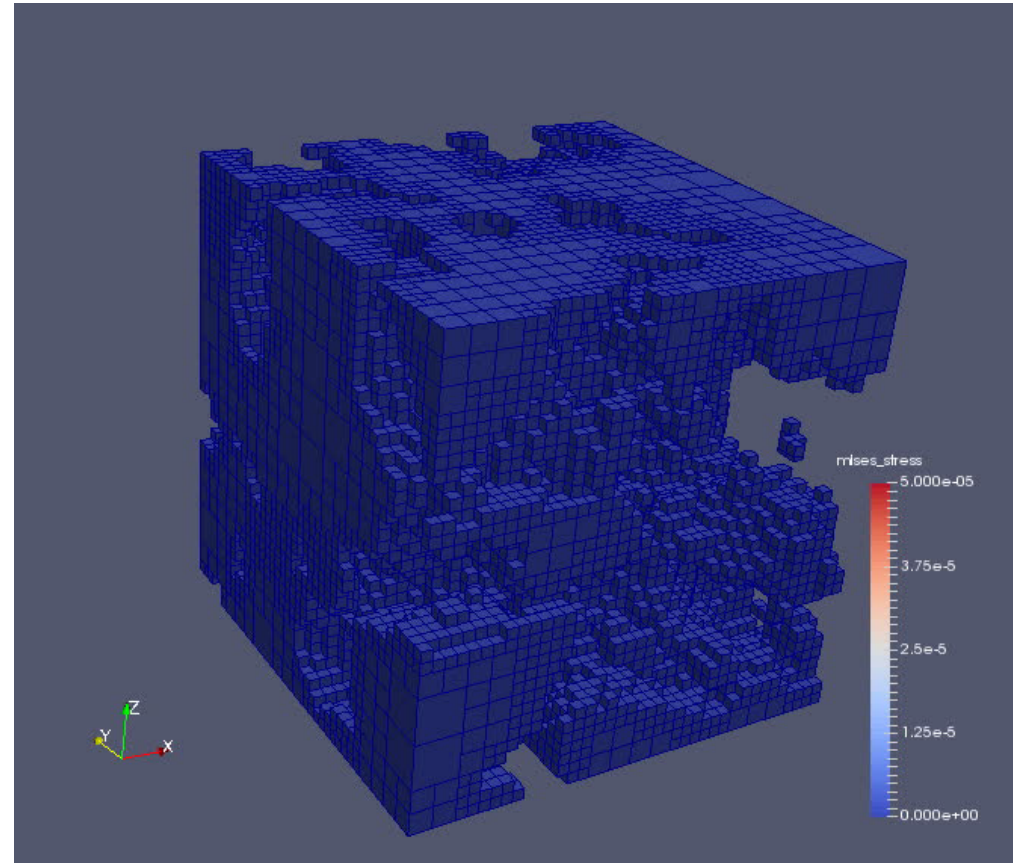
$$r^+ = C_{AB}k^+(T) = [S][F]R_1R_2$$

$$r^- = C_A C_B k^-(T) = 2nf[I][S][F]^2 \frac{R_2^3 R_3}{(R_4^+ - R_4^-)}$$

How does it work?



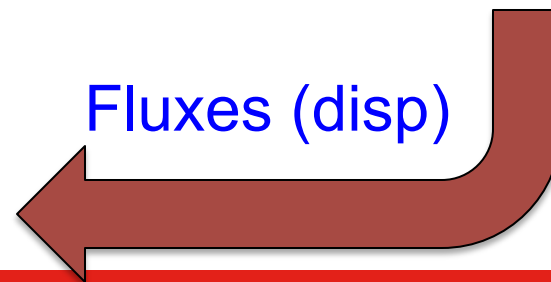
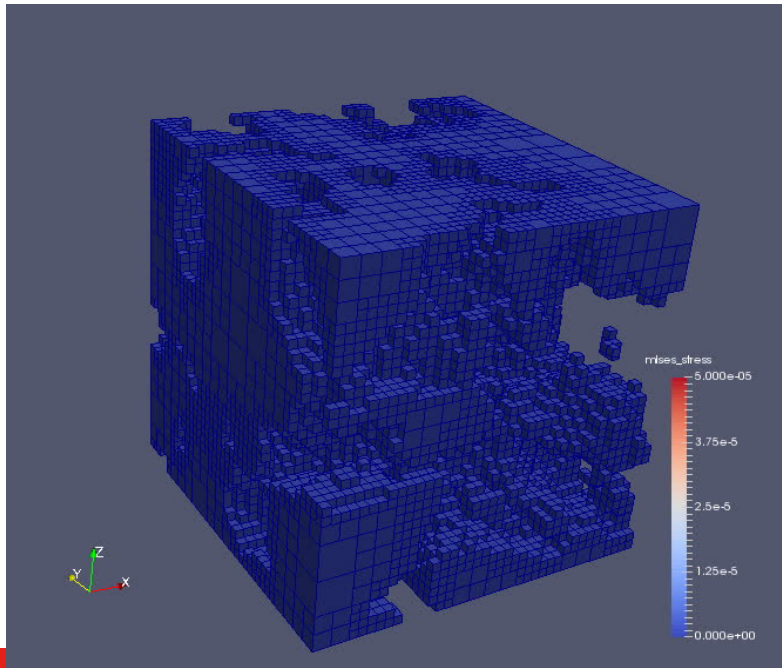
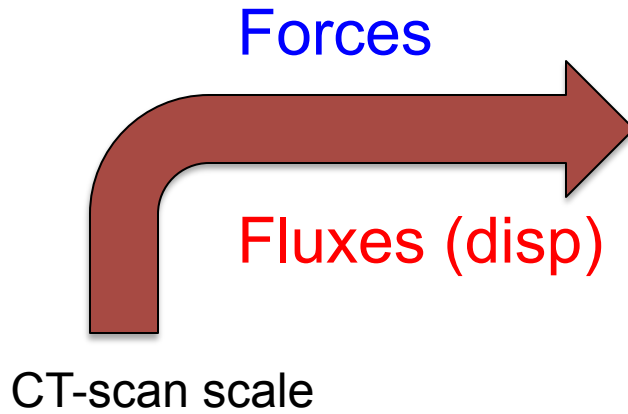
Navier-Stokes fluid flow (hydraulic homogenisation)



See also the poster by Martin Lesueur

How does it work?

Laboratory scale
THMC



Navier-Stokes eqs
Linear elasto-plasticity

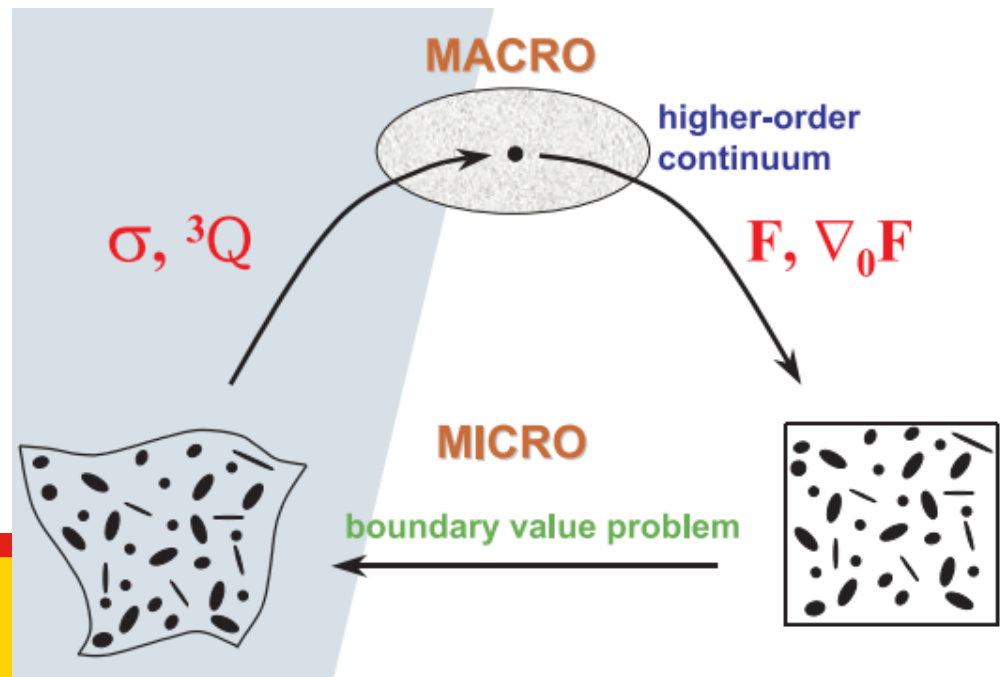
Forces



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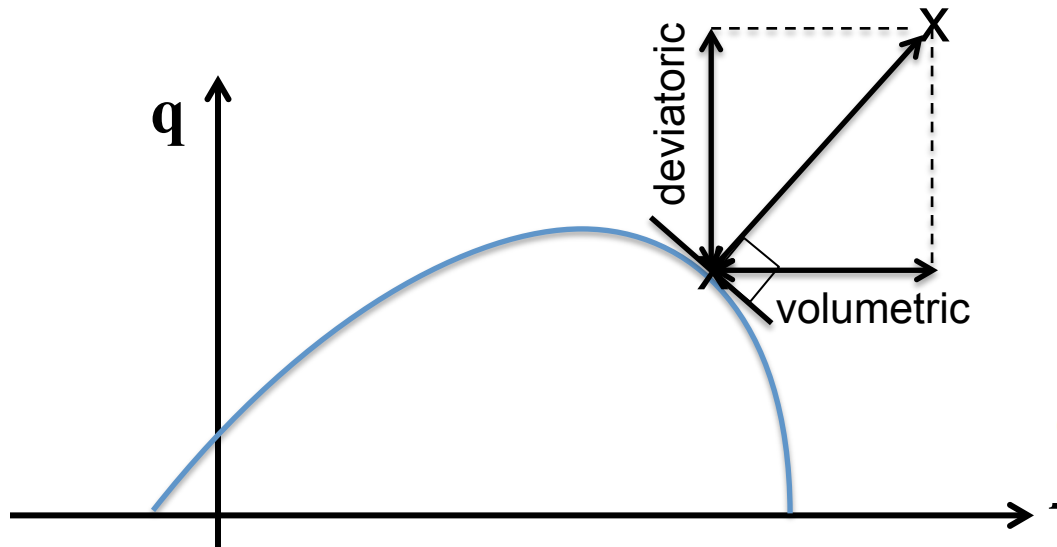
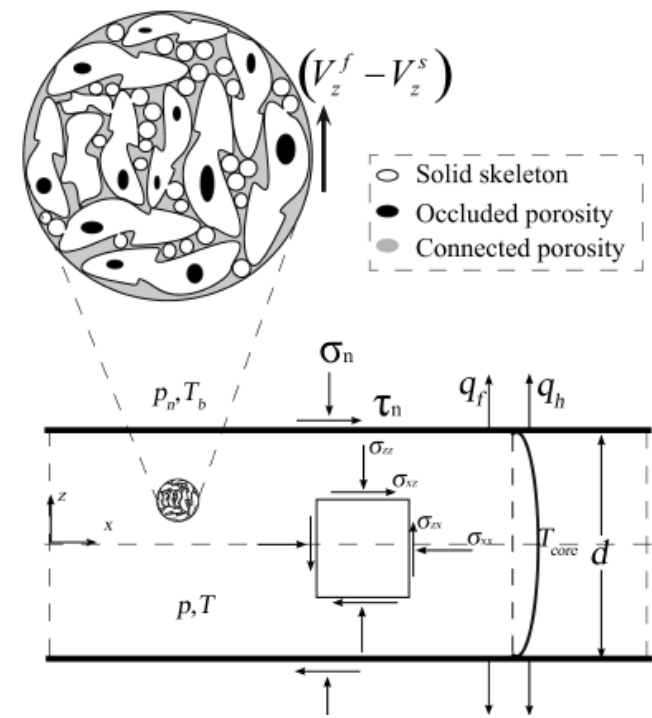
How do we scale up from lab scale to geodynamics?

- THMC systems obey MaxEP, MinEP or something else? Does the choice of BCs make any difference?
- Current paradigm in computational homogenisation follows the MaxEP path
- Is MaxEP representative for THMC systems? How do we deal with the cascade of length/time scales of localisation?



The multiphysics (THMC) approach

- Fluid-saturated rock
- Coaxial Elasto-visco-plasticity, deviatoric and volumetric components



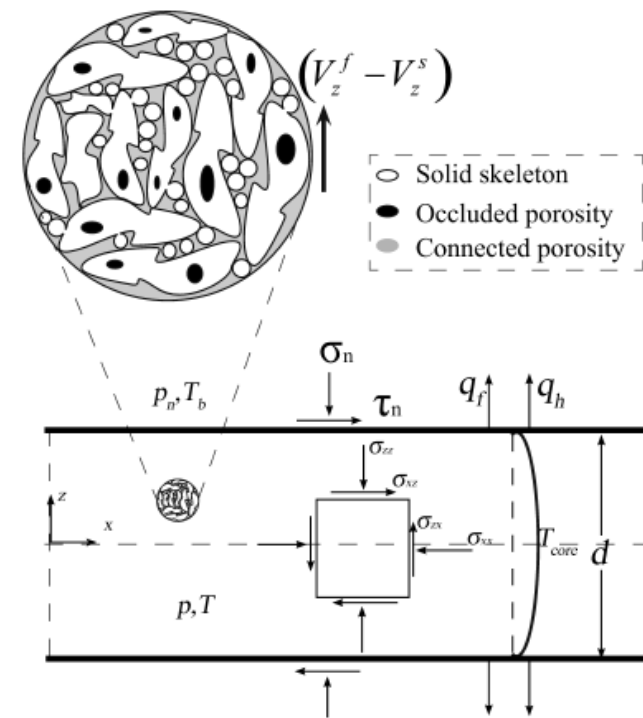
$$\dot{\epsilon}_{ij}^i = \lambda \frac{\partial f}{\partial \sigma'_{ij}} \quad \lambda = \sqrt{\dot{\epsilon}_d^i{}^2 + \dot{\epsilon}_v^i{}^2}$$

$$\dot{\epsilon}_d^i = \dot{\epsilon}_0 \left\langle \frac{q - q_Y}{\sigma_{ref}} \right\rangle^m \exp\left(-\frac{Q_{mech}^d}{RT}\right)$$

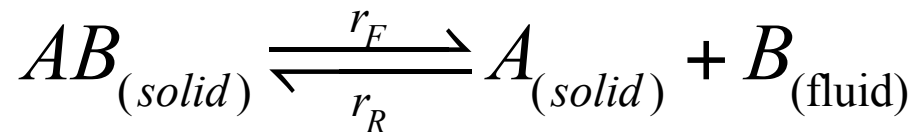
$$\dot{\epsilon}_v^i = \dot{\epsilon}_0 \left\langle \frac{p' - p_Y}{\sigma_{ref}} \right\rangle^m \exp\left(-\frac{Q_{mech}^v}{RT}\right)$$

Towards a unified THMC approach

- Fluid-saturated rock
- Coaxial Elasto-visco-plasticity, deviatoric and volumetric components
- Mechanical (Shear) heating
- Endothermic fluid release reaction producing excess pore pressure



$$\phi = \phi_0 + \Delta\phi_{mech} + \Delta\phi_{chem}$$



- Porosity and permeability linked with Kozeny-Carman law

$$k_{\pi} = k_{\pi 0} \frac{(1 - \phi_0)^2}{\phi_0^3} \frac{\phi^3}{(1 - \phi)^2}$$

$$\Delta\phi_{chem} = A_{\phi} \frac{1 - \phi_0}{1 + \frac{\rho_B}{\rho_A} \frac{M_A}{M_B} \frac{1}{s}},$$

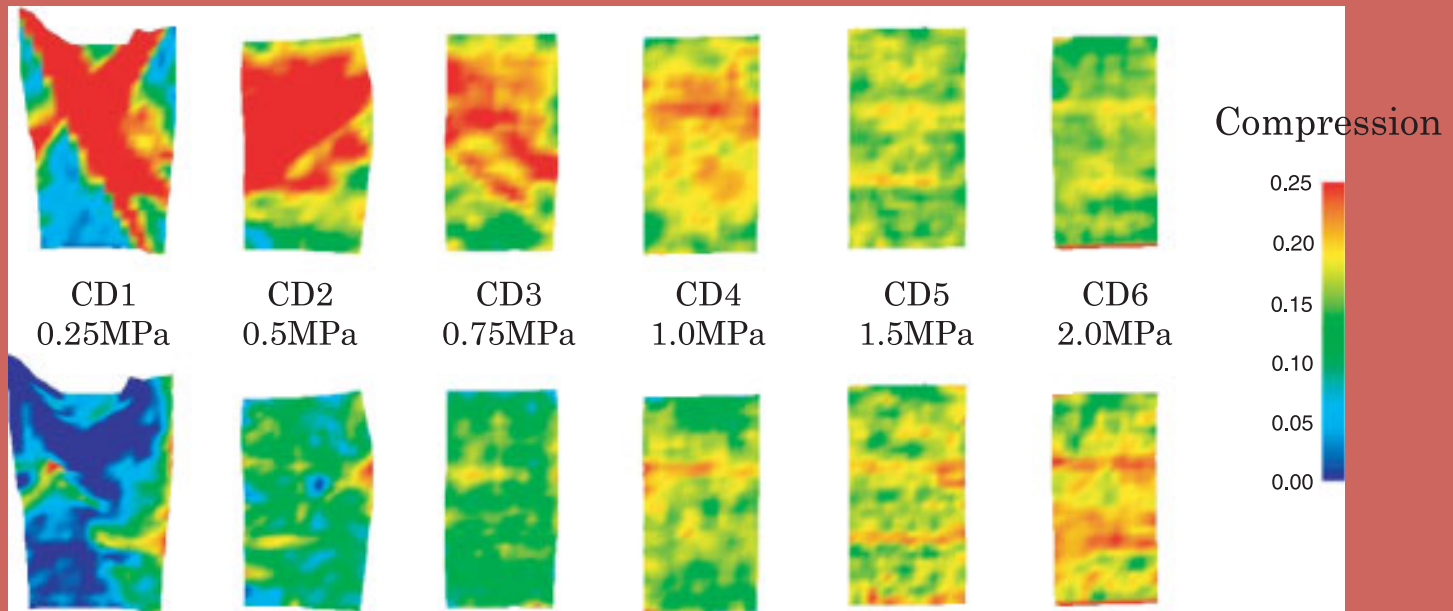
$$s = \frac{\omega_{rel}}{1 + \omega_{rel}}, \text{ and}$$

$$\omega_{rel} = \frac{\rho_{AB}}{\rho_A} \frac{M_A}{M_{AB}} K_c \exp\left(\frac{\Delta H}{RT}\right)$$

Do we need chemistry?

Triaxial experiments (THM) in soft rocks

Distribution of shear strain



Distribution of volumetric strain

Oka et al, IJNAMG, 2011

The goal is to
constrain the
hardening law:

$$Q_{mech} = E + p_f V_{act}$$

$$E = E_0 + \Delta E$$

a (T)HM system for pore collapse

$$0 = \partial_j \sigma'_{ij} - \partial_i \Delta p_f + b_i,$$

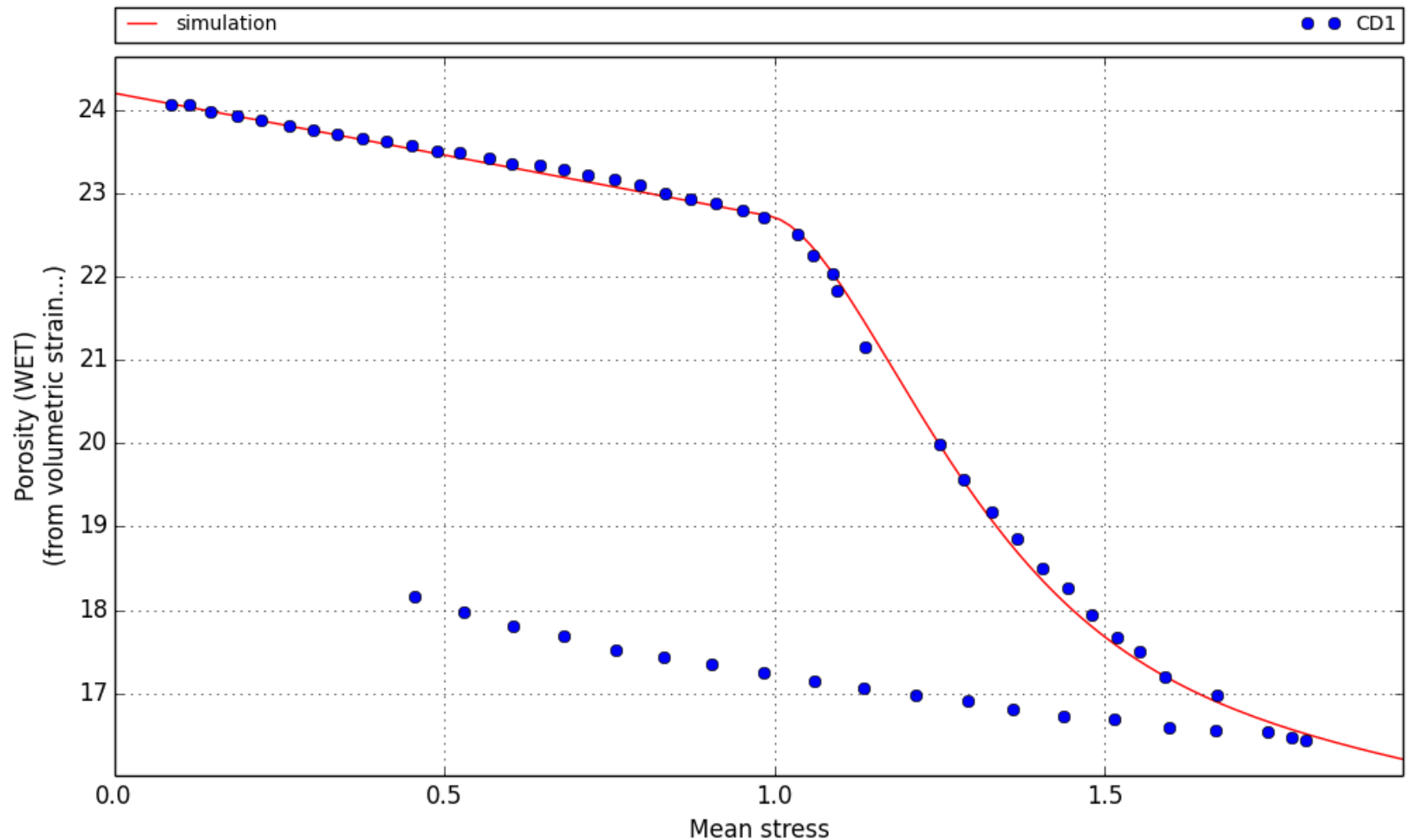
$$0 = \partial_t \Delta p_f - \partial_i \left[\frac{1}{Le} \partial_i \Delta p_f \right] - \Lambda \partial_t T + \frac{\dot{\epsilon}_V}{\beta},$$

$$0 = \partial_t T - \partial_{ii}^2 T - Gr \sigma_{ij} \dot{\epsilon}_{ij}^{pl}.$$

The energy and entropy balance laws are solved explicitly to determine the global attractor of the system

Diagenesis (pore collapse) in sandstone

Isotropic consolidation of saturated Bleurwiller sandstone

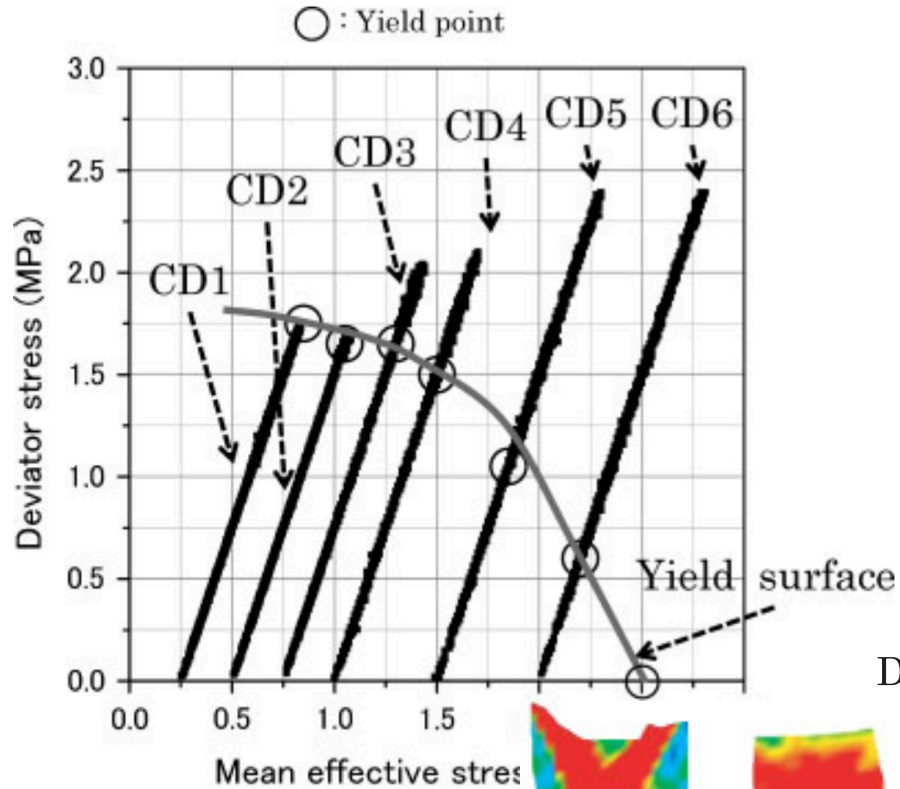


Fortin et al, JGR 2007

$$\Delta f_{\text{nonint}}^{\text{pores}} = p \frac{3(1 - \nu_o)}{4E_o} \left\{ \frac{10(1 + \nu_o)}{7 - 5\nu_o} \text{tr}(\boldsymbol{\sigma} \cdot \boldsymbol{\sigma}) - \left[\frac{1 + 5\nu_o}{7 - 5\nu_o} + \frac{1}{3(1 + \delta_s)} \right] (\text{tr } \boldsymbol{\sigma})^2 \right\},$$

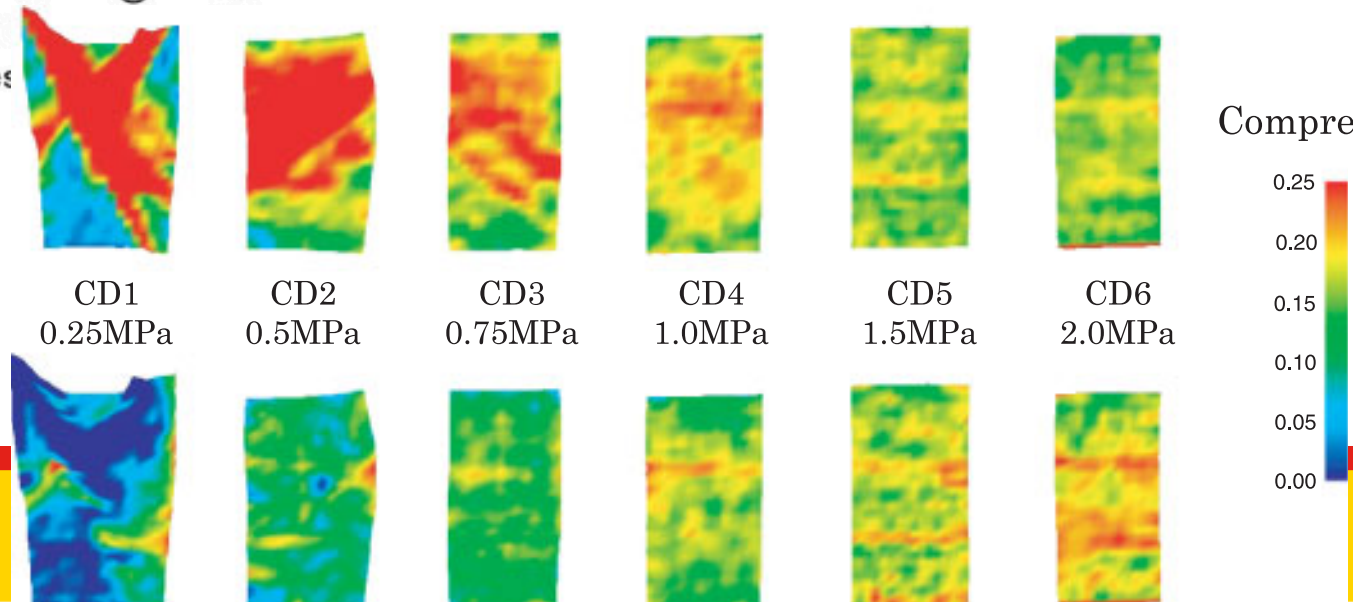


Drained consolidation of diatomaceous mudstone

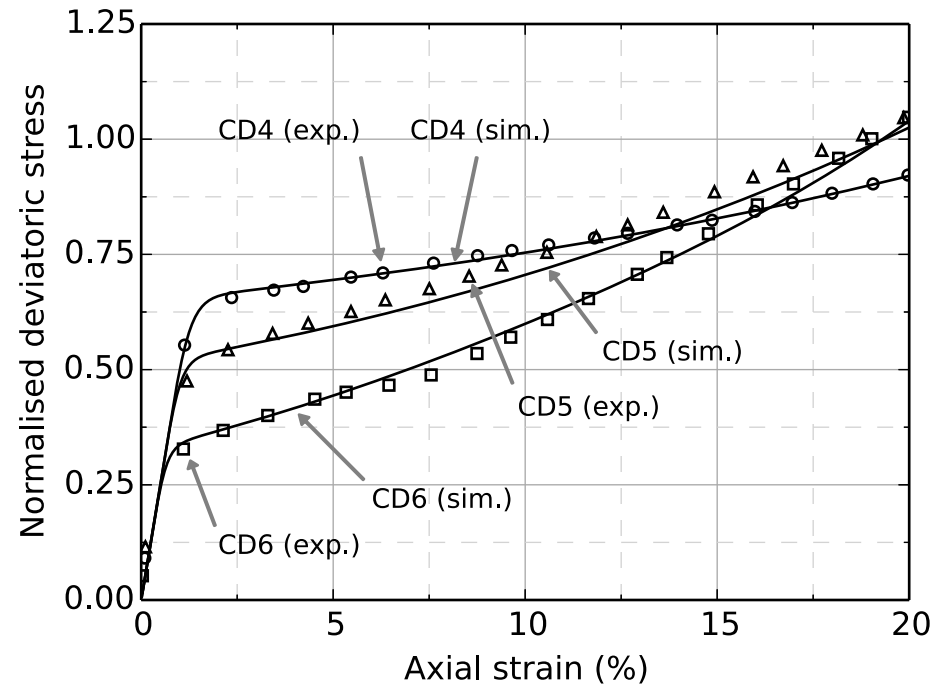
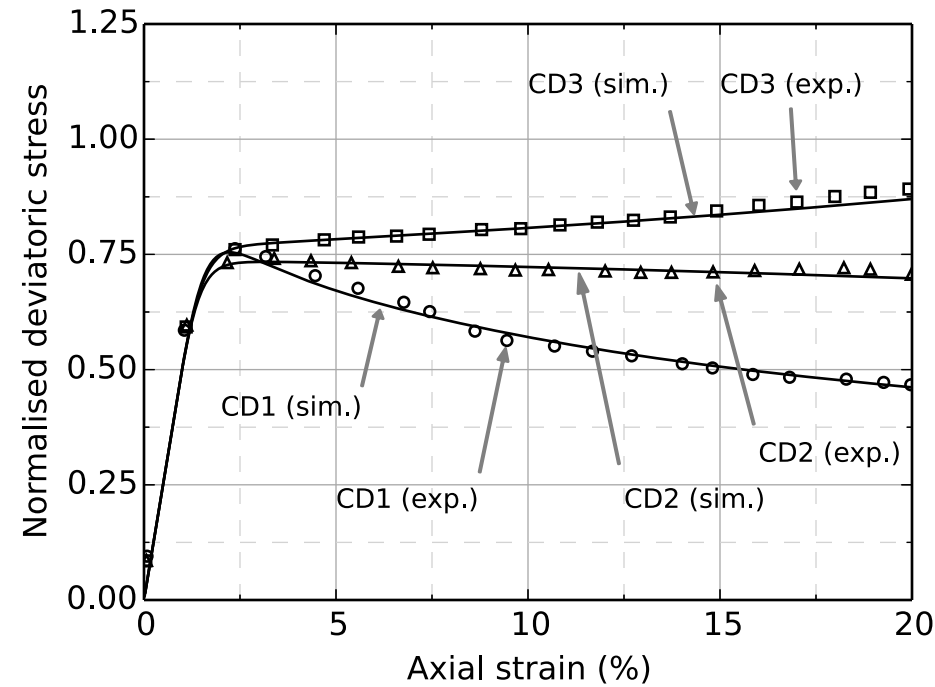


Oka et al, IJNAMG, 2011

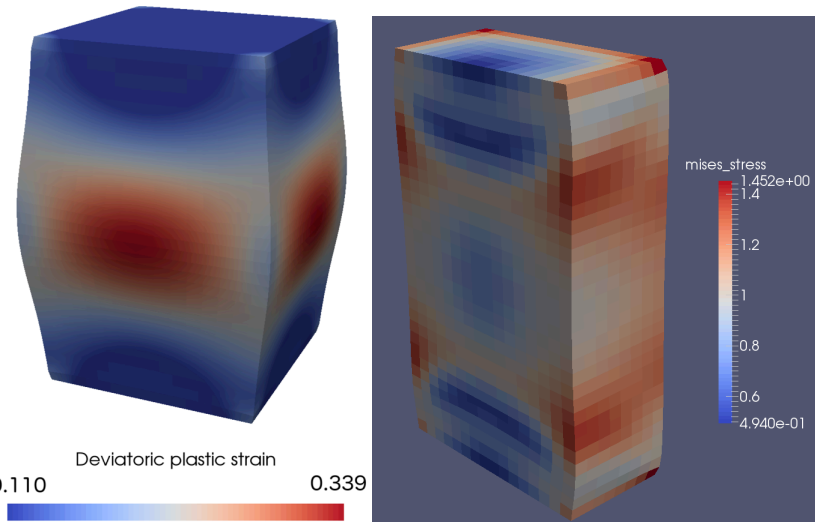
Distribution of shear strain



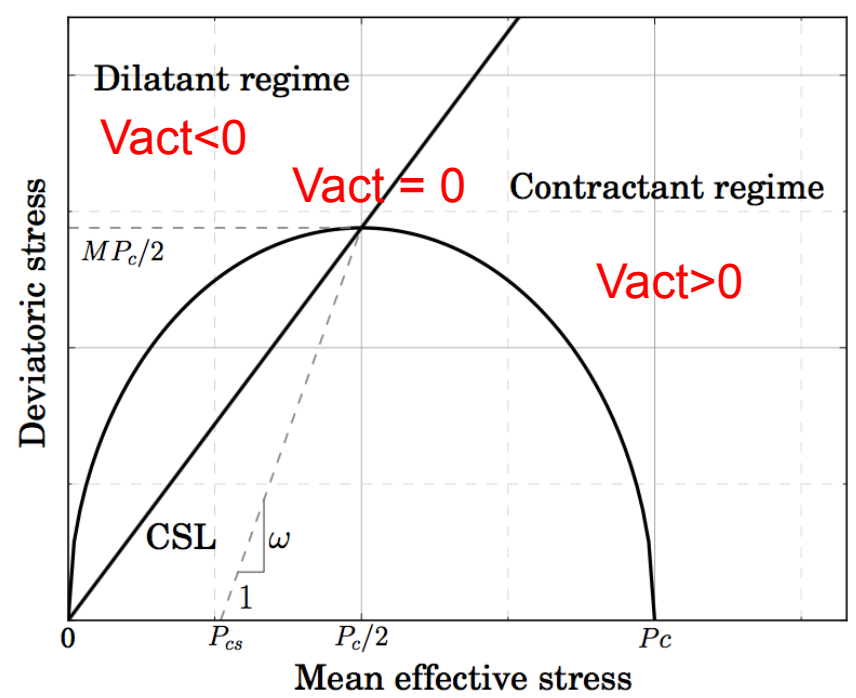
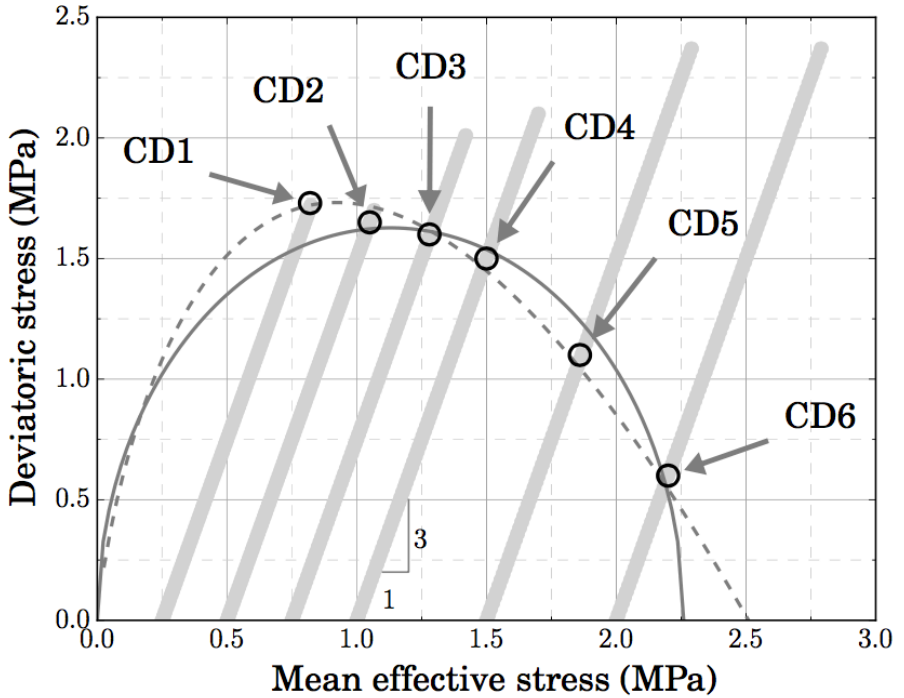
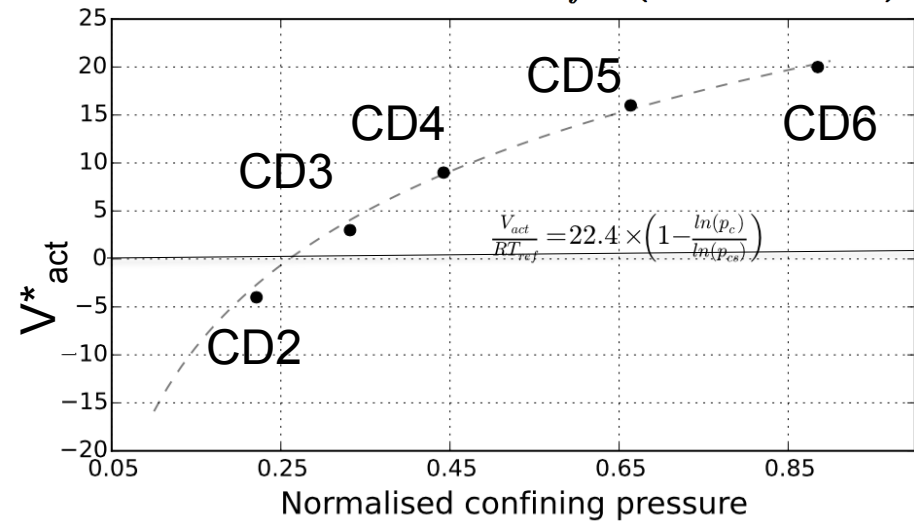
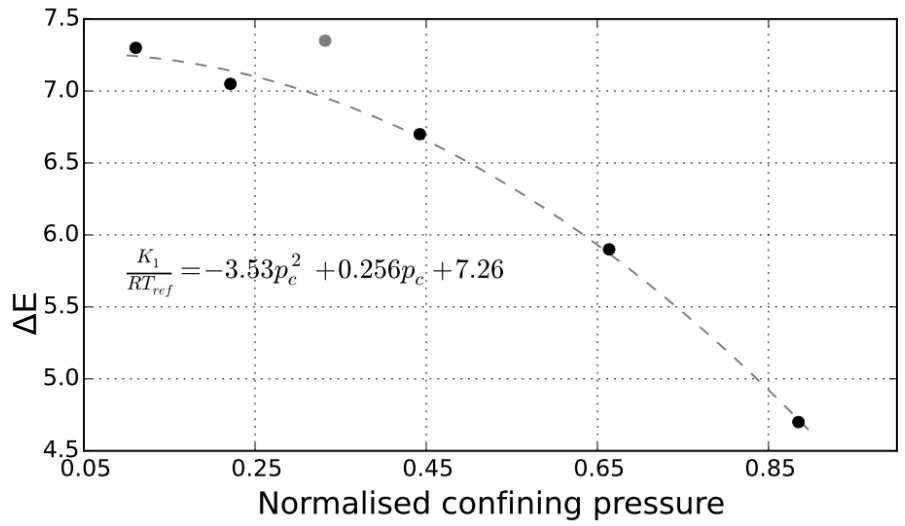
Matching the experiments



See also poster by Mustafa Sari for more materials (sandstone, carbonates...)

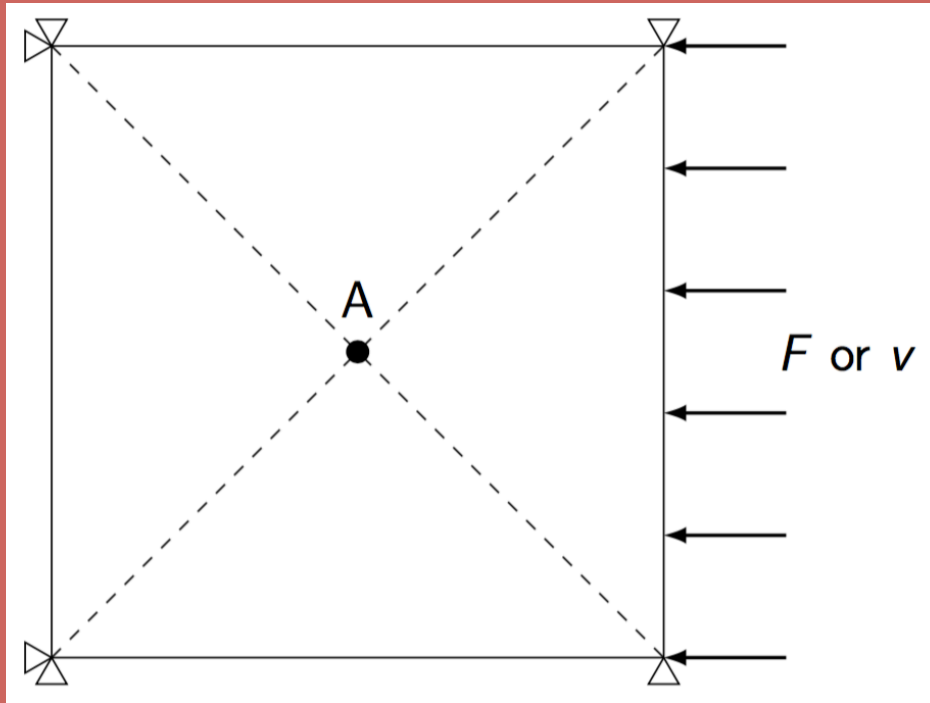


Retrieving information for the rheology $V_{act}^{(w)} = \frac{h_{act}^{(w)}}{\sigma_{ref}} \left(1 - \frac{\ln P_c}{\ln P_{cs}} \right)$

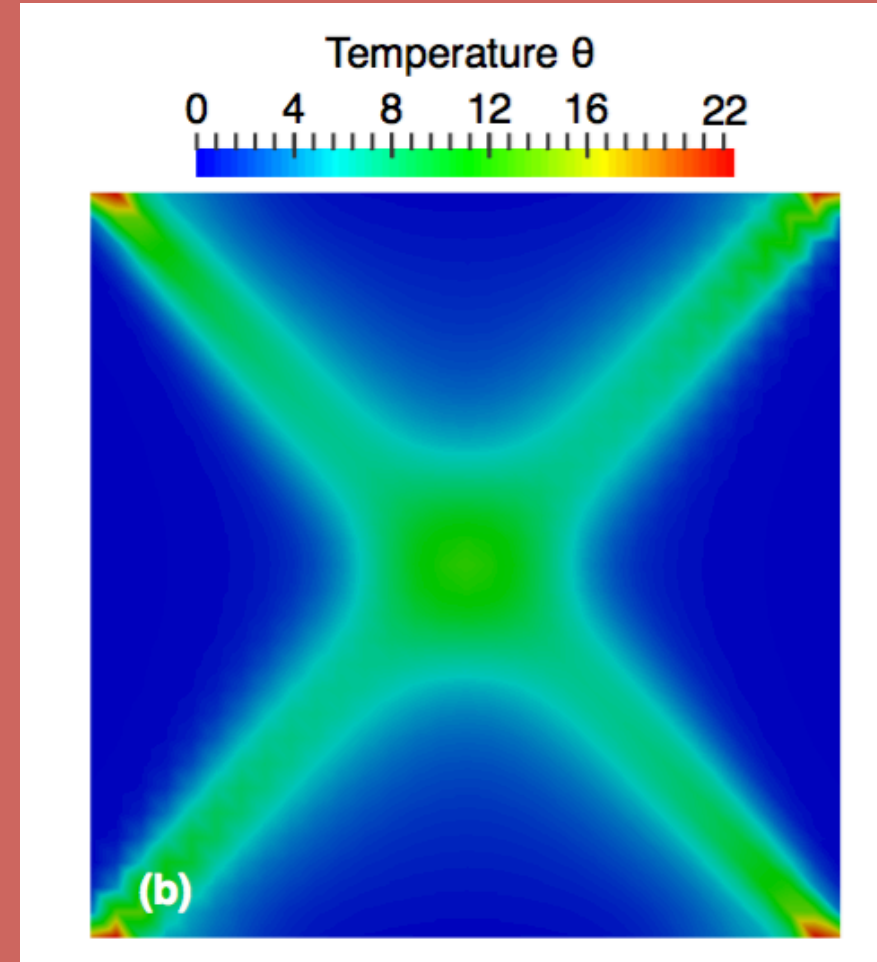


MaxEP or MinEP

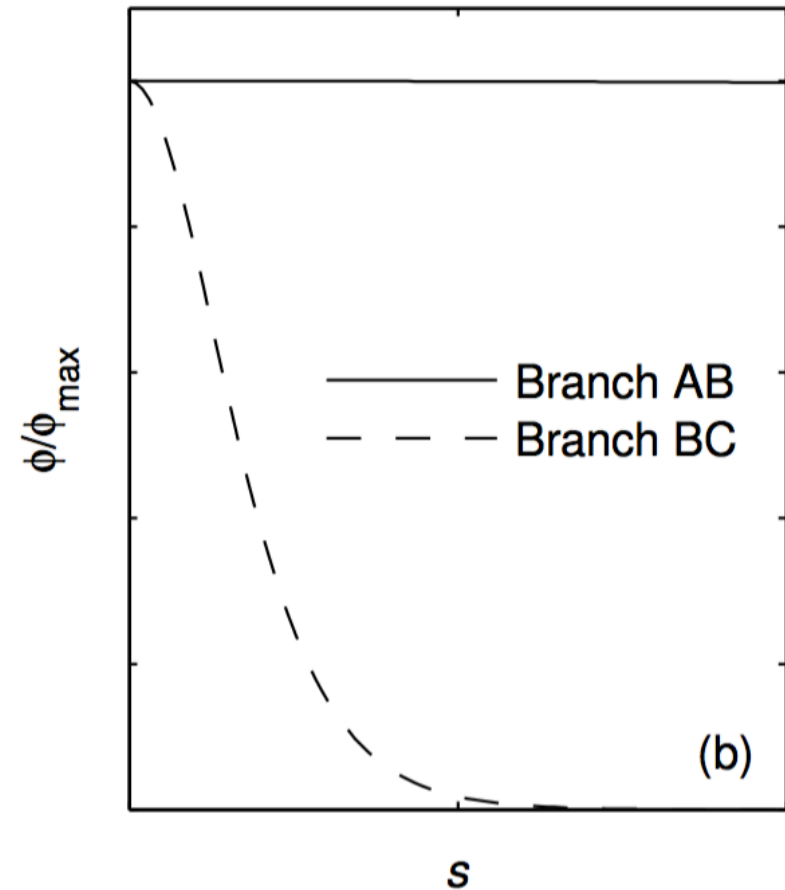
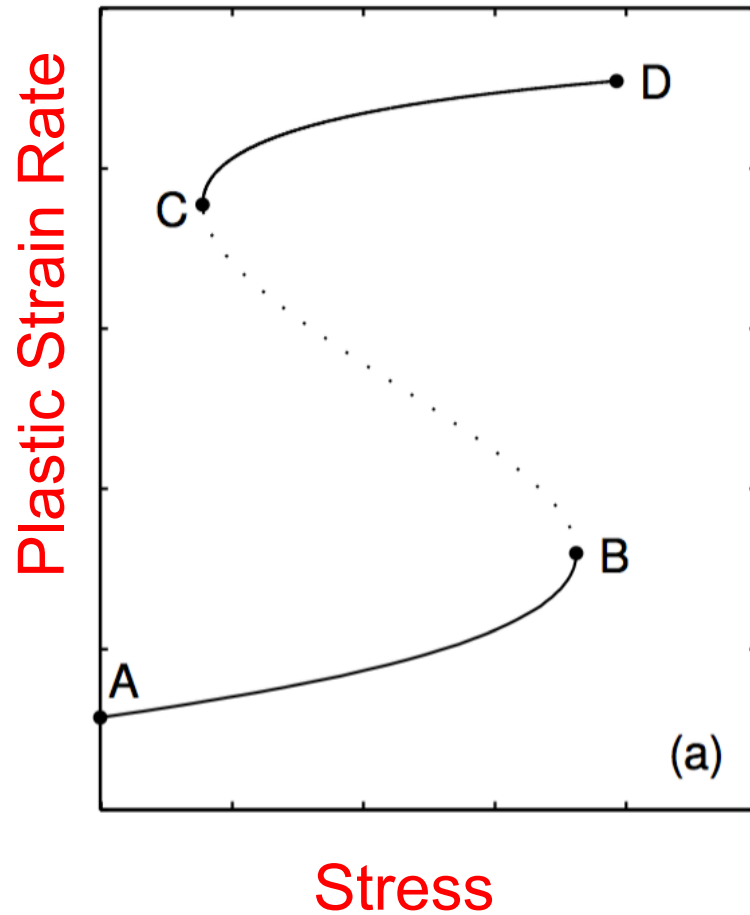
Where do THM systems self-organise?



Paesold et al JoMMS 2016

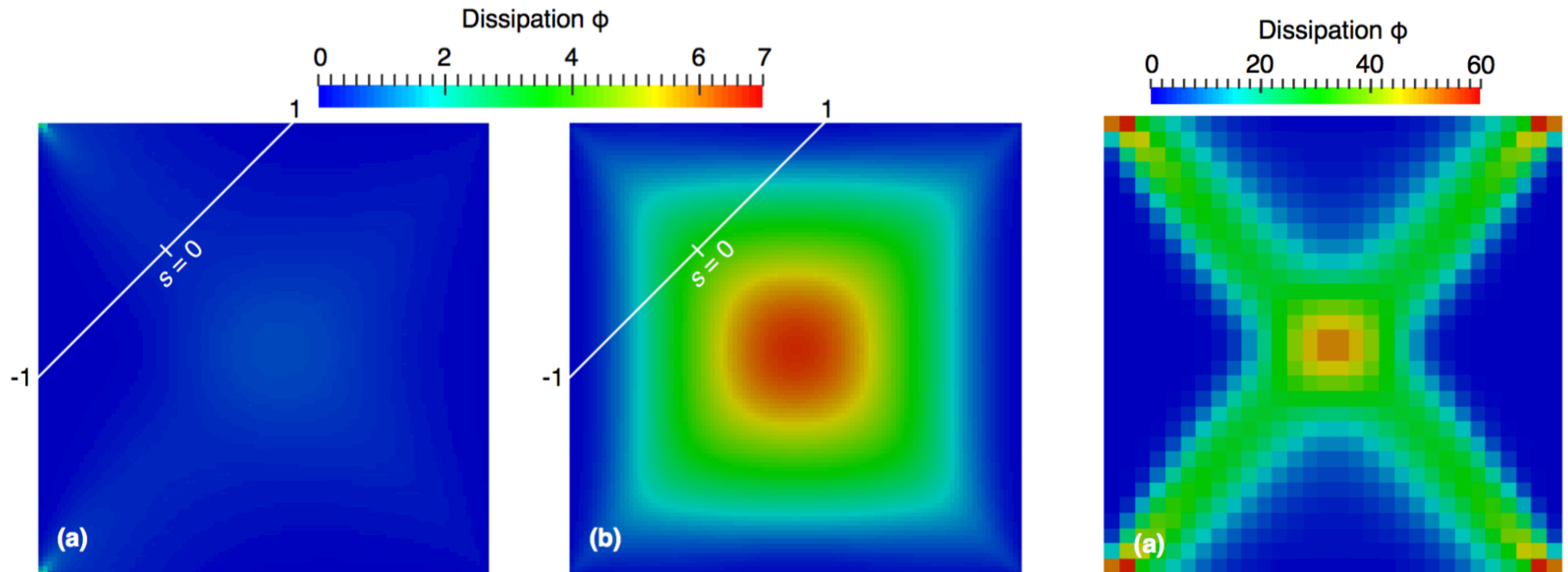


T(H)M systems self-organise at MaxEP



Paesold et al JoMMS 2016

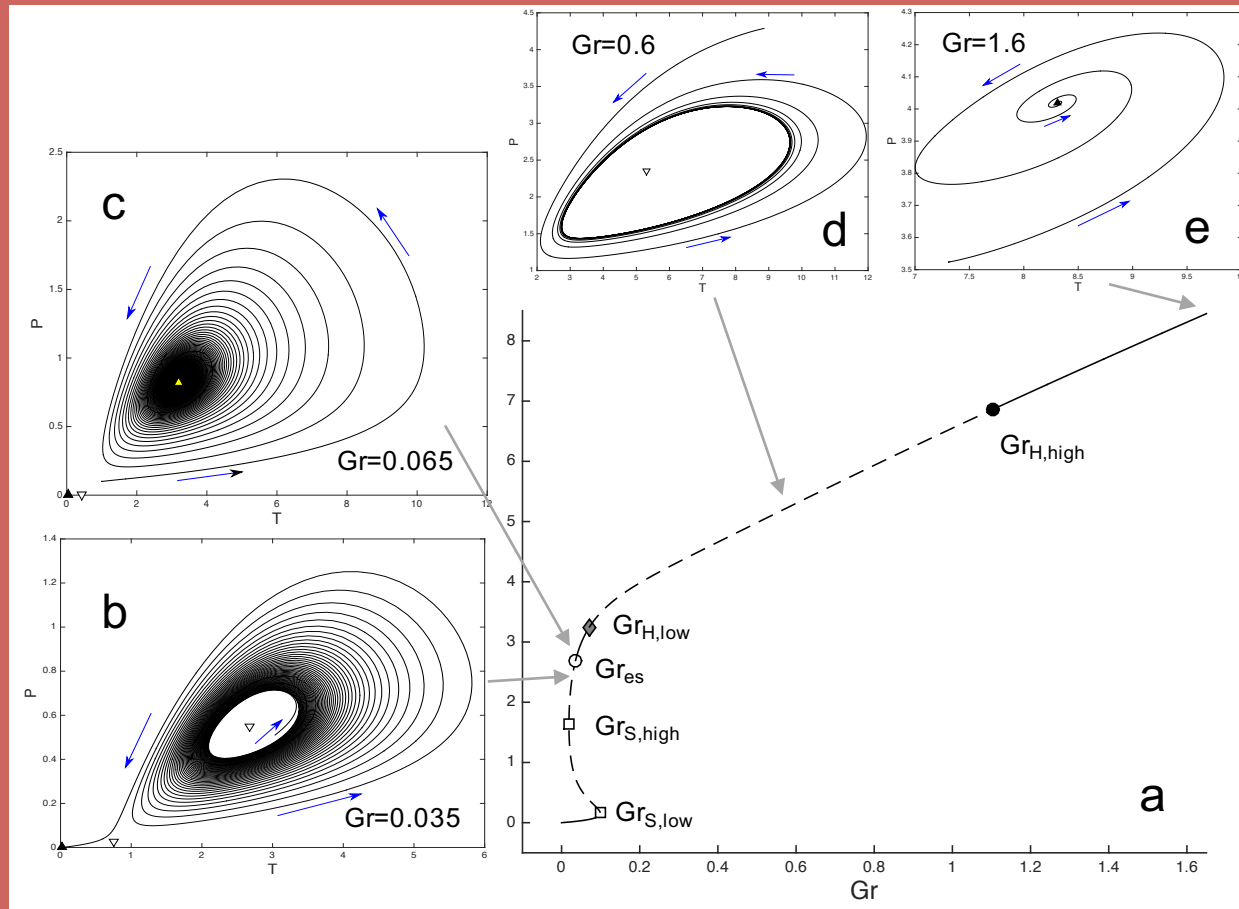
T(H)M systems self-organise at MaxEP



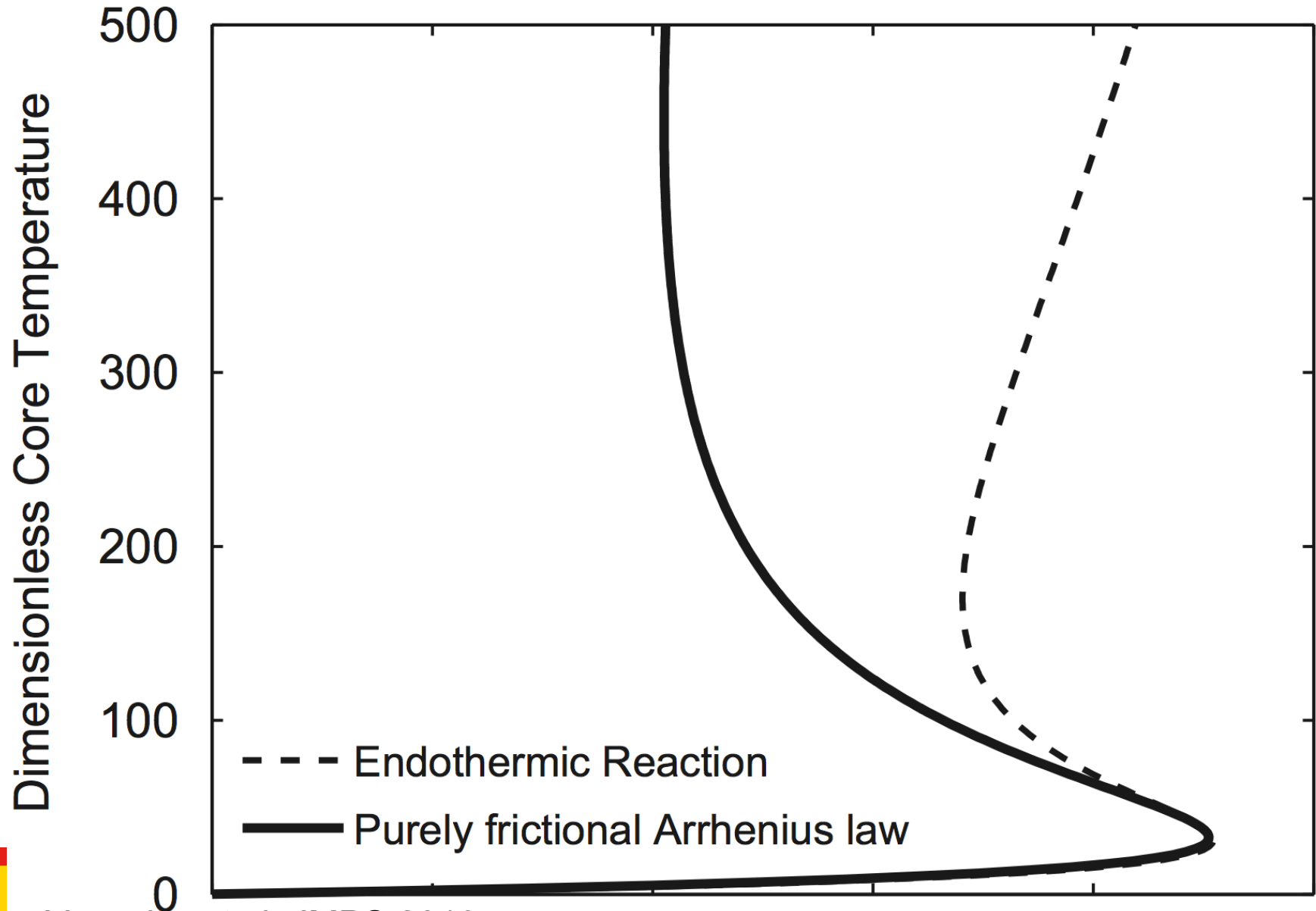
Paesold et al JoMMS 2016

So...

Where do THMC systems self-organise?



What changes when chemistry is included?



System of equations

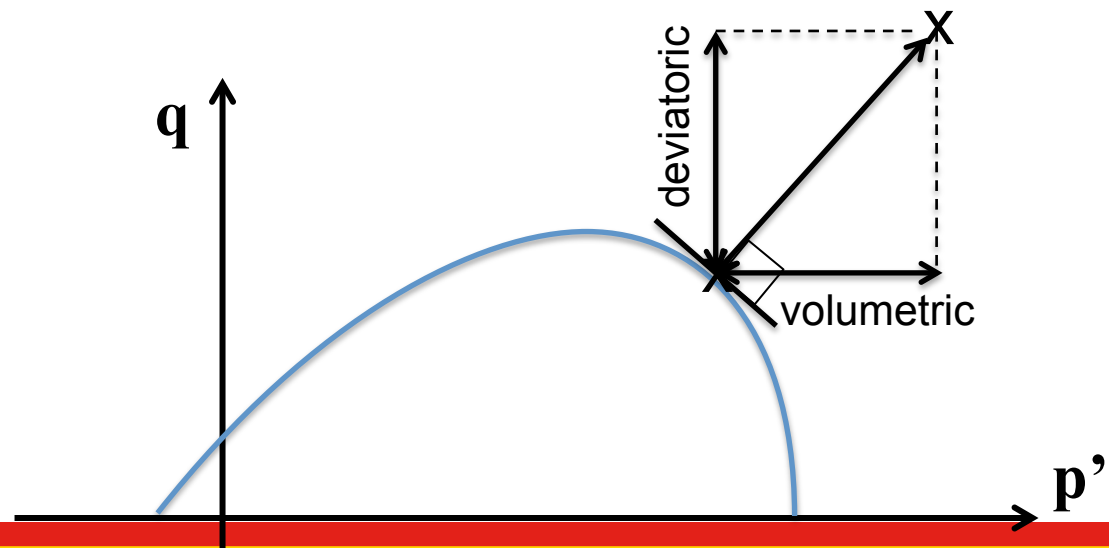
$$AB_{(solid)} \xrightleftharpoons[r_R]{r_F} A_{(solid)} + B_{(fluid)} \quad r_{F,R} = k_{F,R} \rho_i e^{-\frac{Ar}{1+\theta}}$$

$$\frac{\partial P}{\partial t} + Pe_{mass} \mathbf{v} \cdot \nabla P - Pe_{temp} \mathbf{v} \cdot \nabla \theta - \nabla \cdot \left[\frac{1}{Le} \nabla P \right] - \bar{\Lambda} \frac{\partial \theta}{\partial t} + \frac{\dot{\epsilon}_v^{pl}}{\beta^*} - \mu \cdot r_F = 0$$

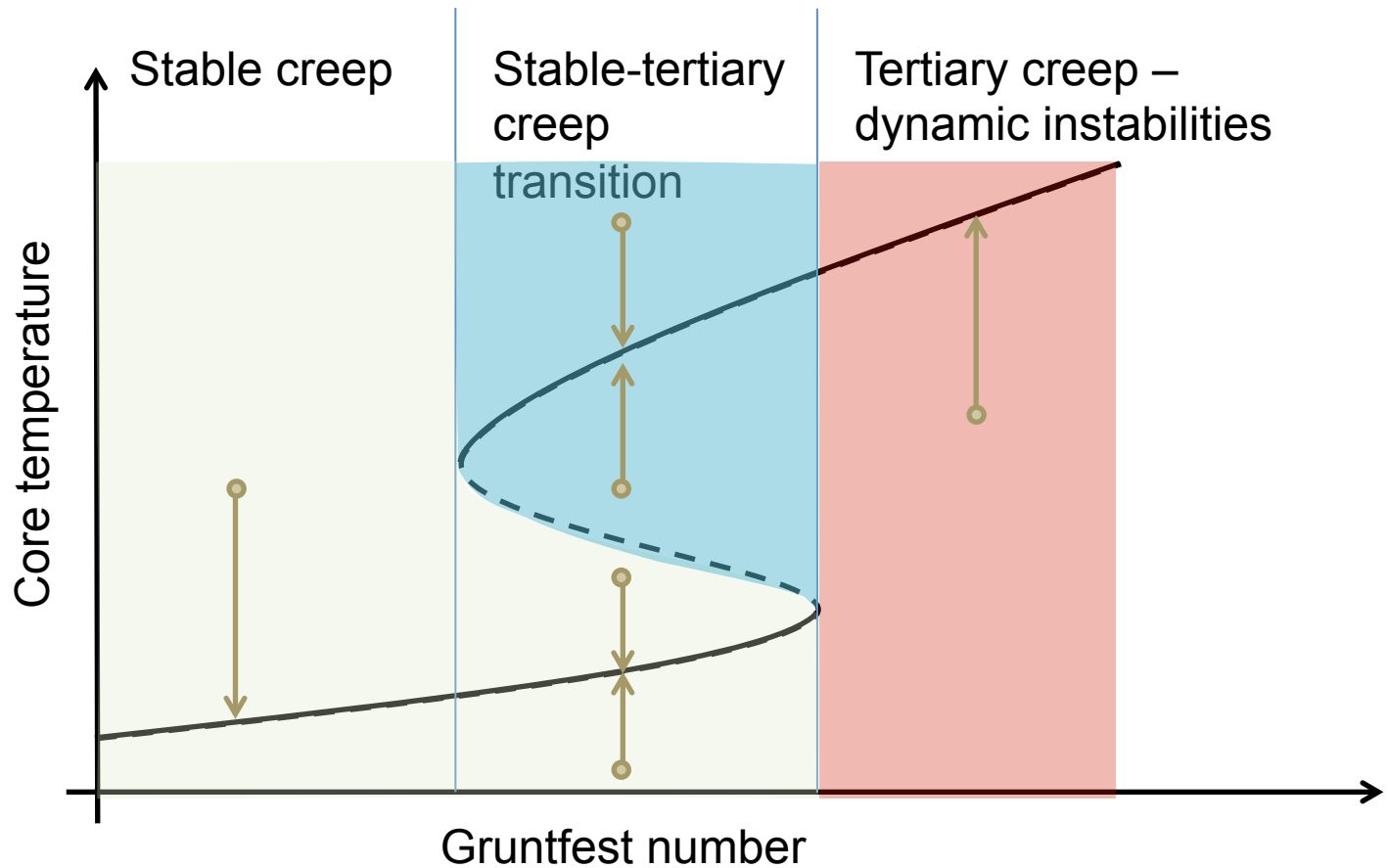
$$\frac{\partial \theta}{\partial t} + Pe_{thermal} \mathbf{v} \cdot \nabla \theta - \nabla^2 \theta - \mathbf{X} \cdot \boldsymbol{\tau} \cdot \dot{\epsilon}_d^p - \mathbf{X} \cdot \boldsymbol{\sigma} \cdot \dot{\epsilon}_v^p + |\Delta H| \cdot r_F - |\Delta H| \cdot r_R = 0$$

$$\frac{\partial(\sigma'_{ij} + P\delta_{ij})}{\partial x_i} = 0$$

$$\dot{\epsilon}_{ij} = C_{ijkl}^{-1} \dot{\sigma}_{kl} + \lambda \frac{\partial f}{\partial \sigma_{ij}}$$



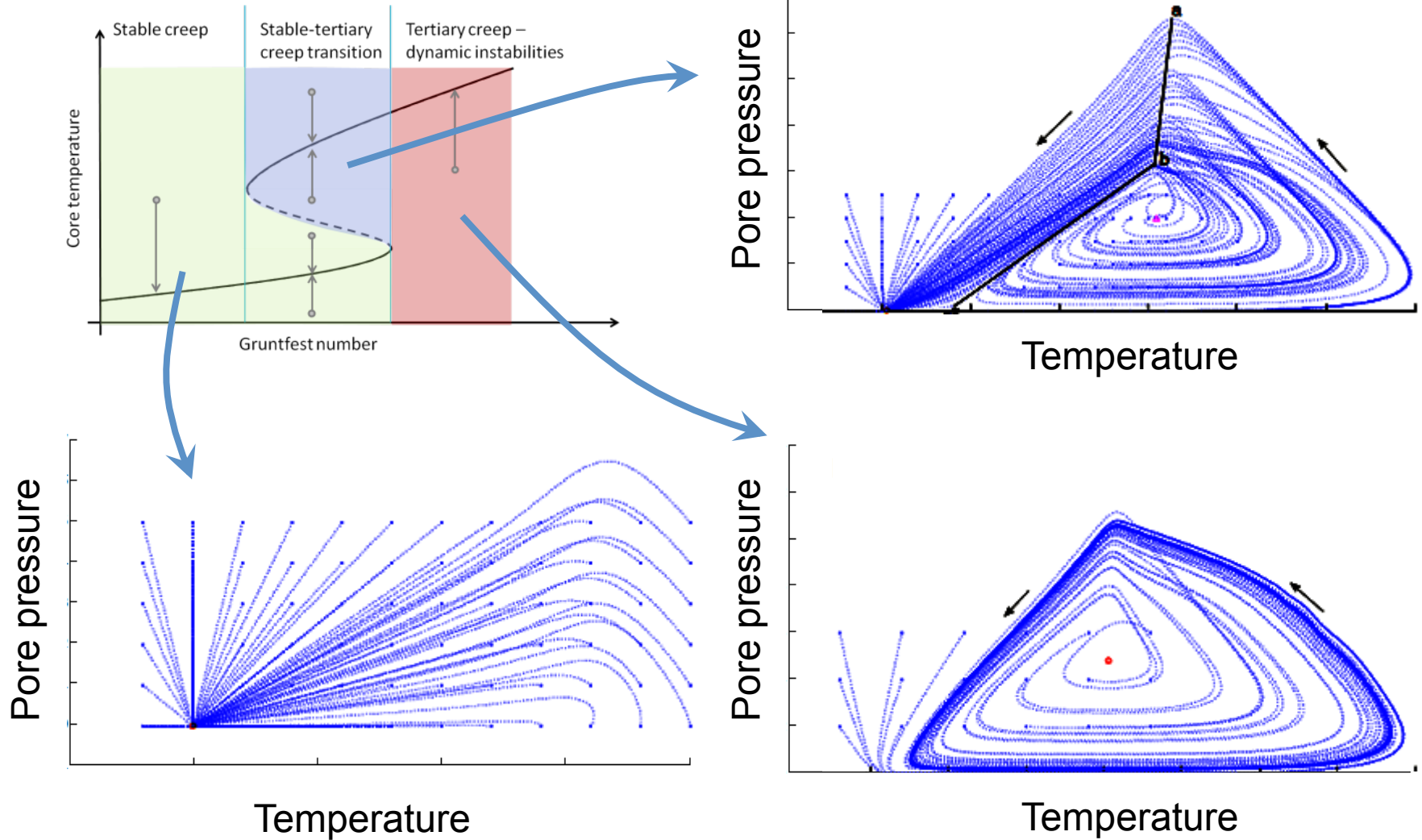
System's stability regimes



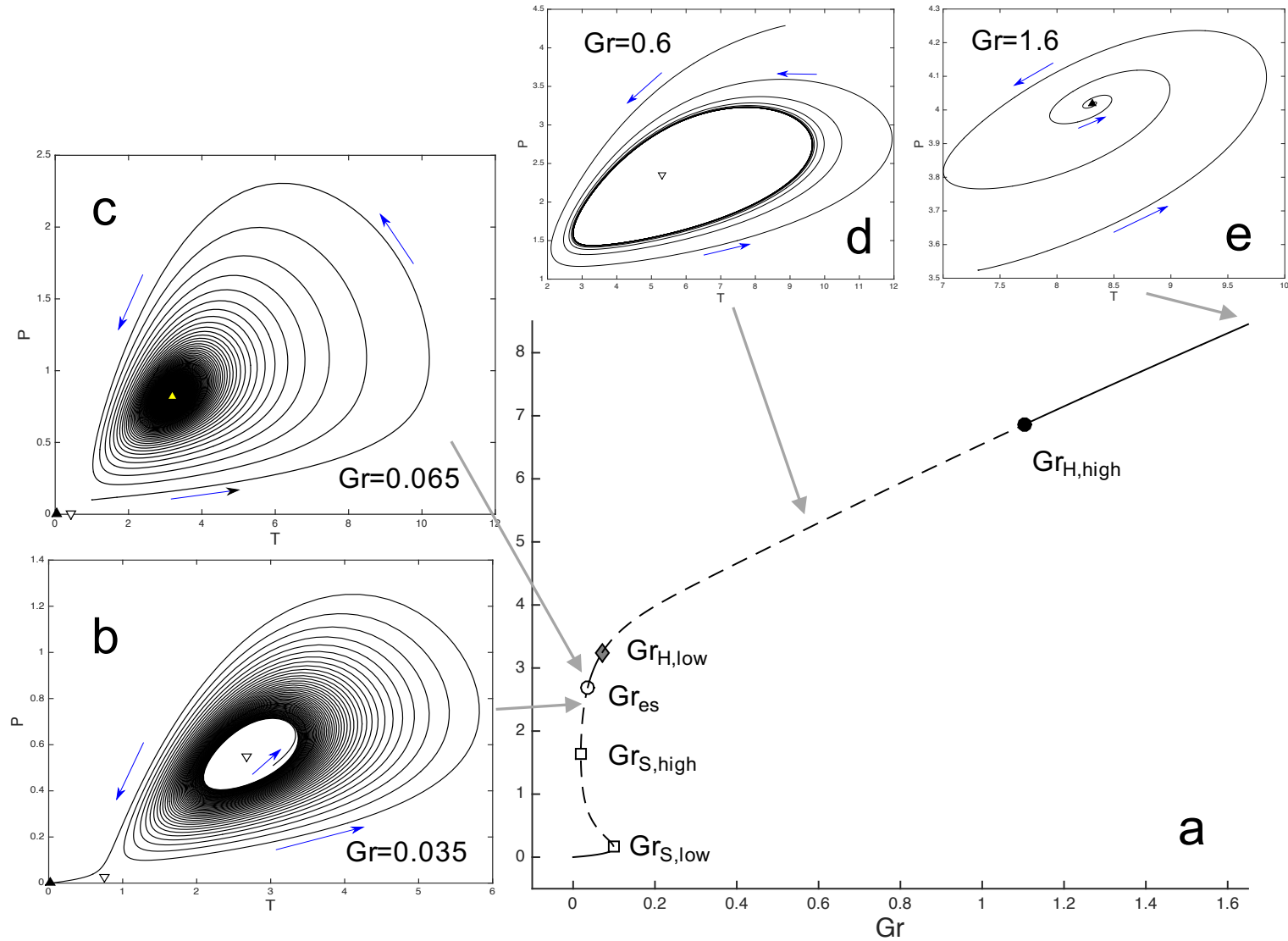
$$Gr = \frac{\text{char. time scale heat production}}{\text{char. time scale energy transfer}}$$

Phase diagrams

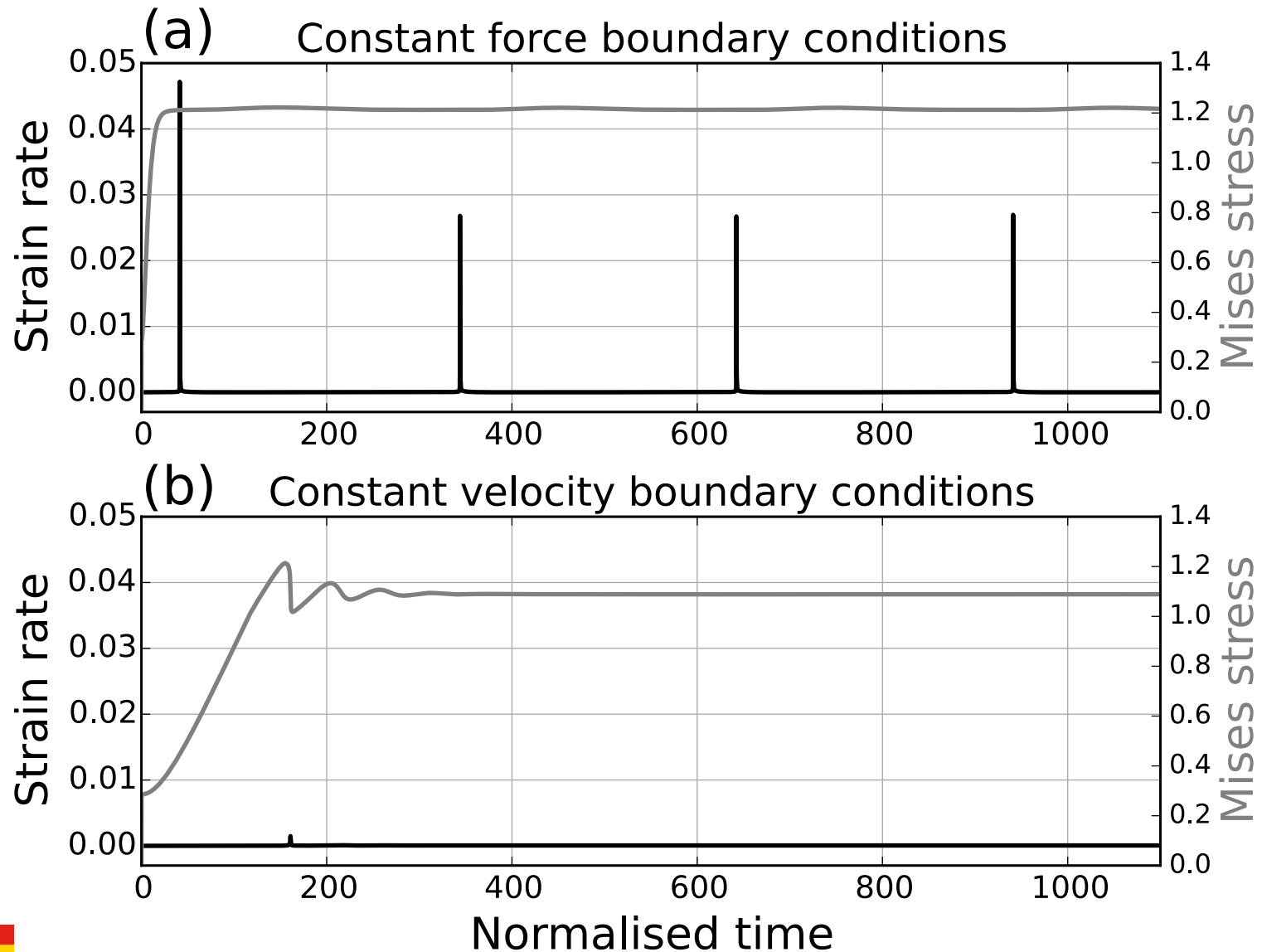
Natural localised instability



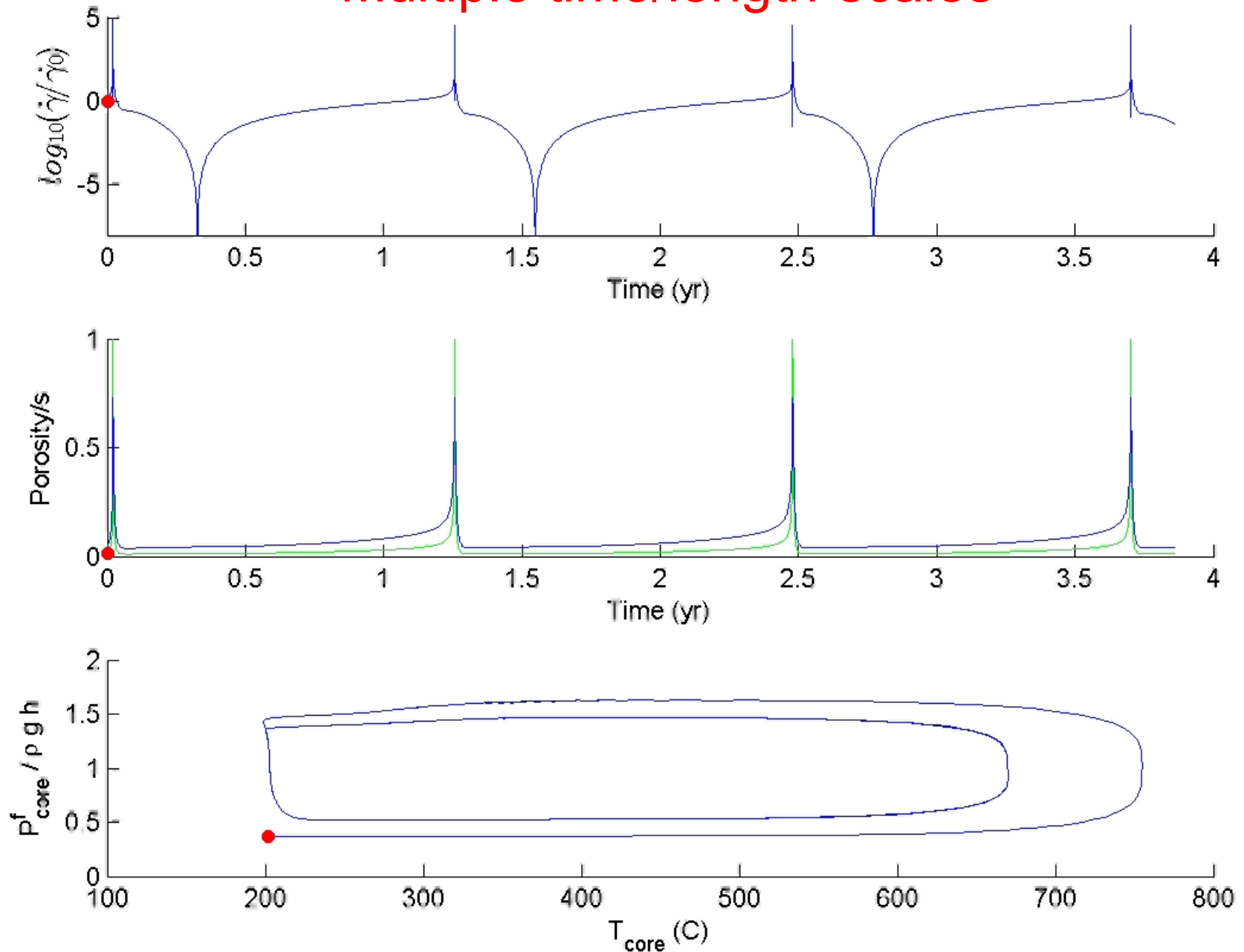
THMC systems self-organise around MinEP



How do we upscale from here? Constant force vs flux BCs



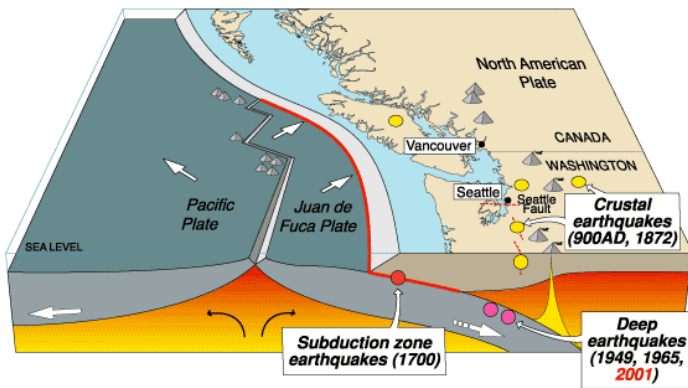
Multiple time/length-scales



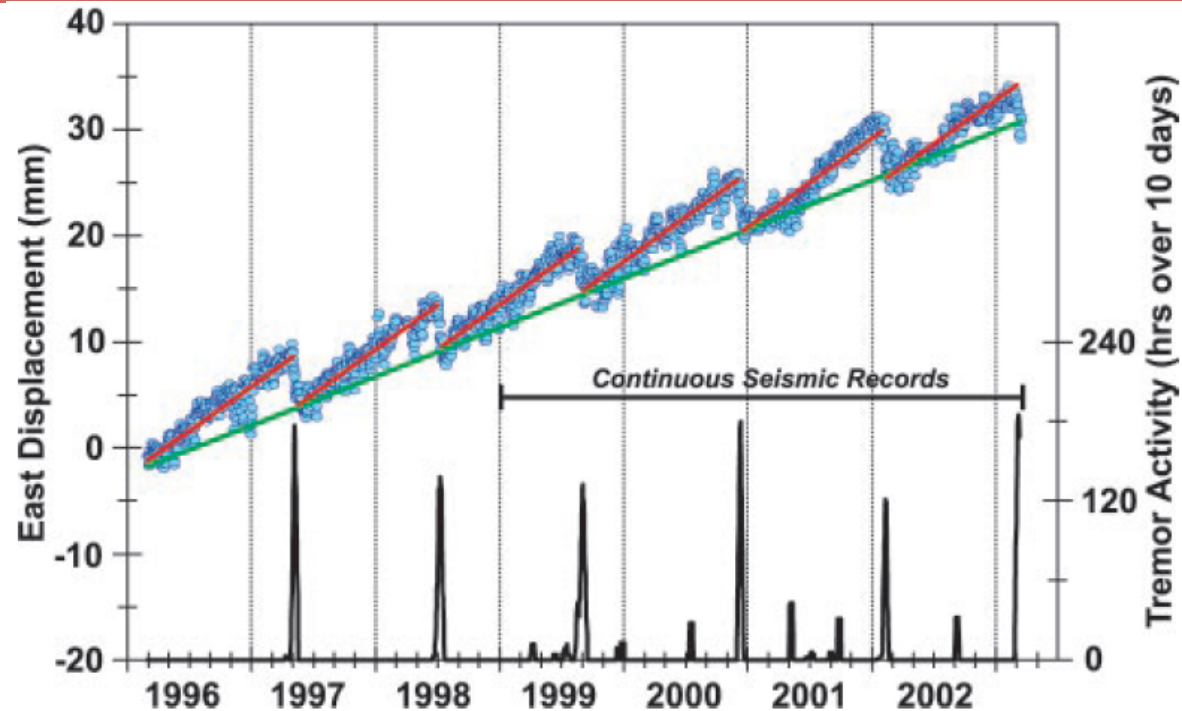
Extrapolation to large scales?

Large scale modelling

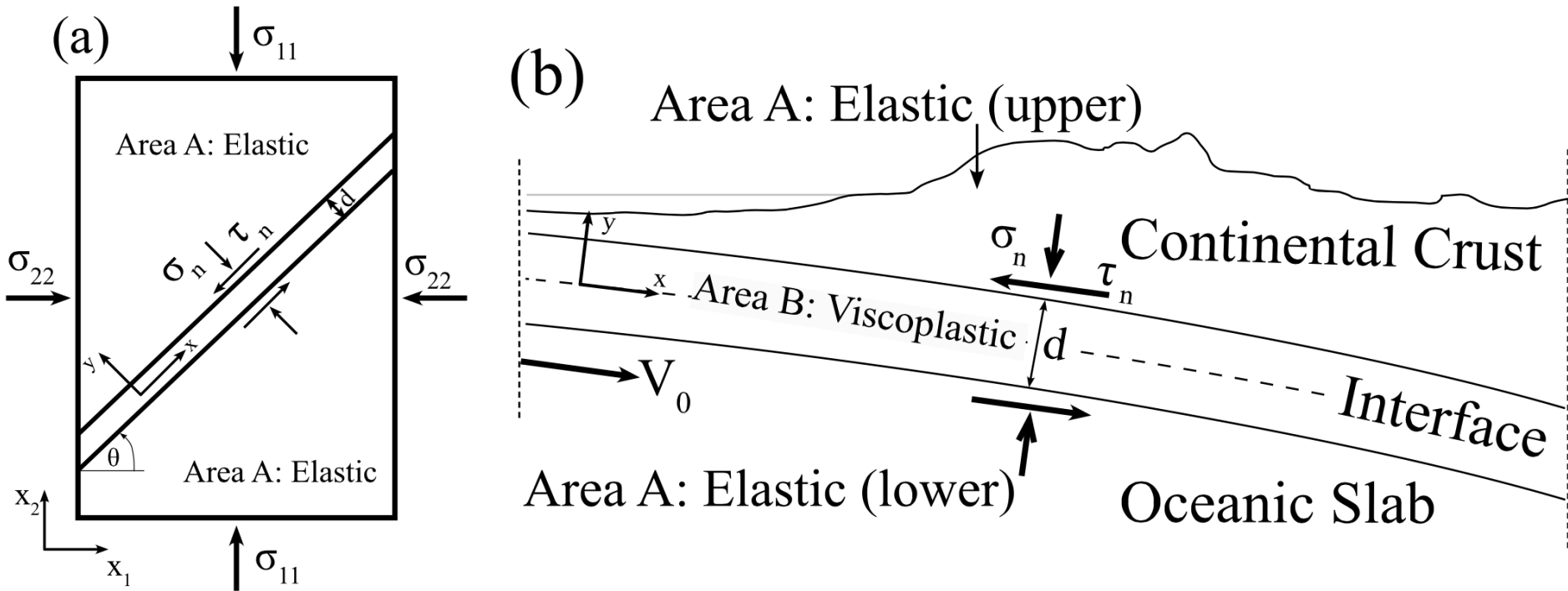
Cascadia earthquake sources



| Source | Affected area | Max. Size | Recurrence |
|---------------------------|---------------|-----------|-----------------|
| ● Subduction Zone | W.WA, OR, CA | M 9 | 500-600 yr |
| ● Deep Juan de Fuca plate | W.WA, OR, | M 7+ | 30-50 yr |
| ● Crustal faults | WA, OR, CA | M 7+ | Hundreds of yr? |

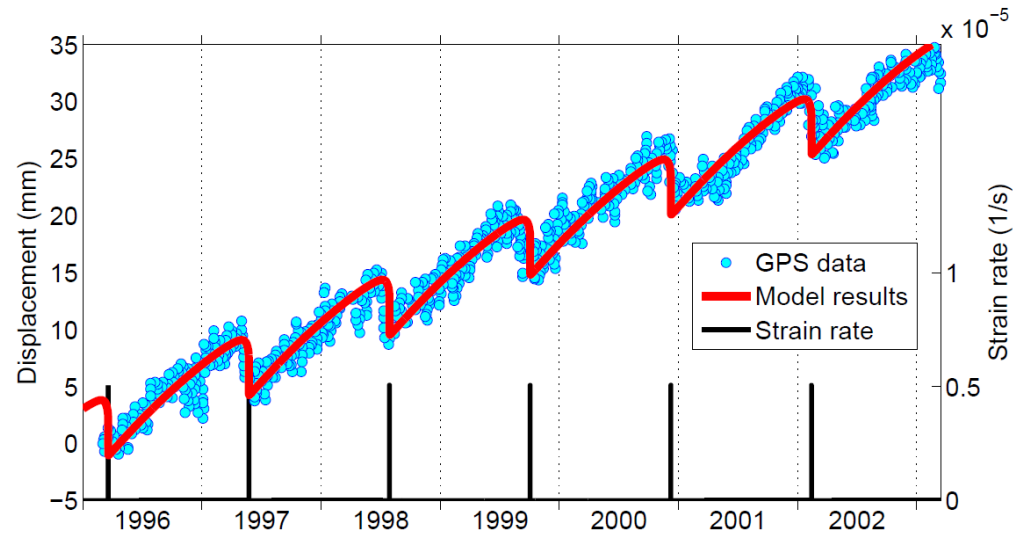
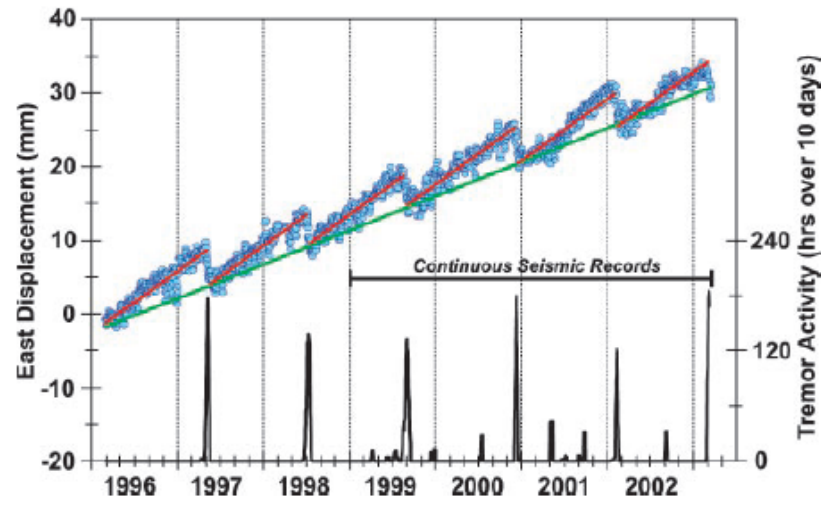
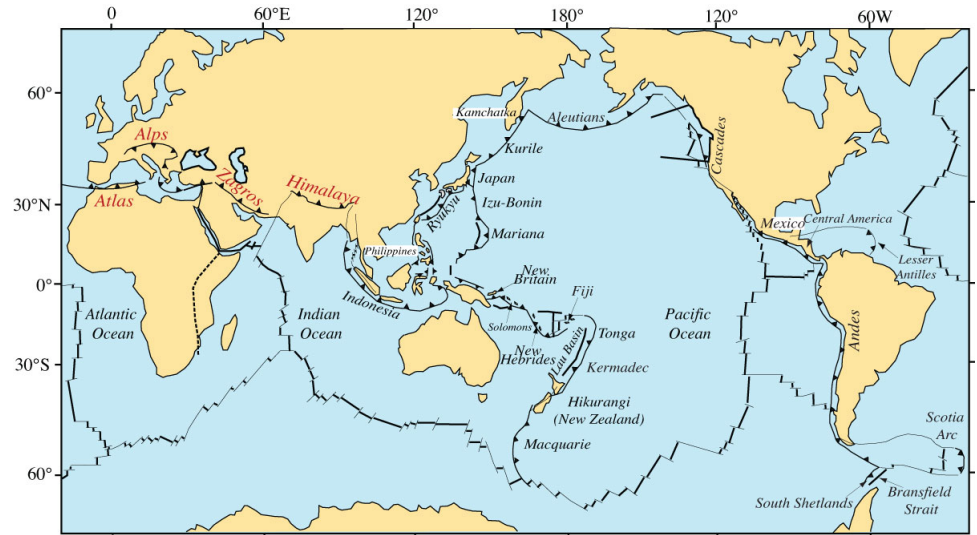
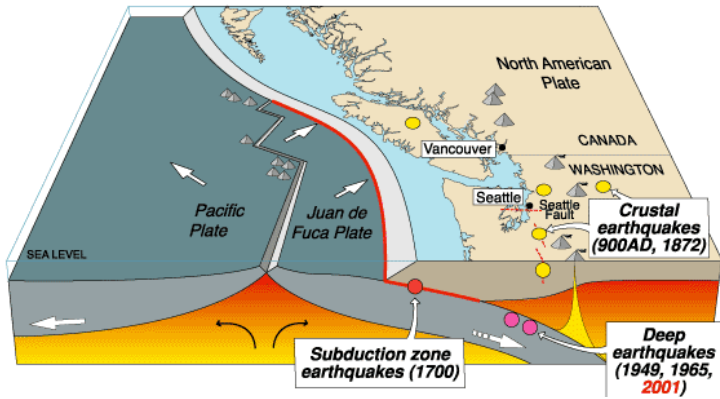


Modelling Subduction zones: Serpentinite dehydration oscillator



Modelling subduction zones

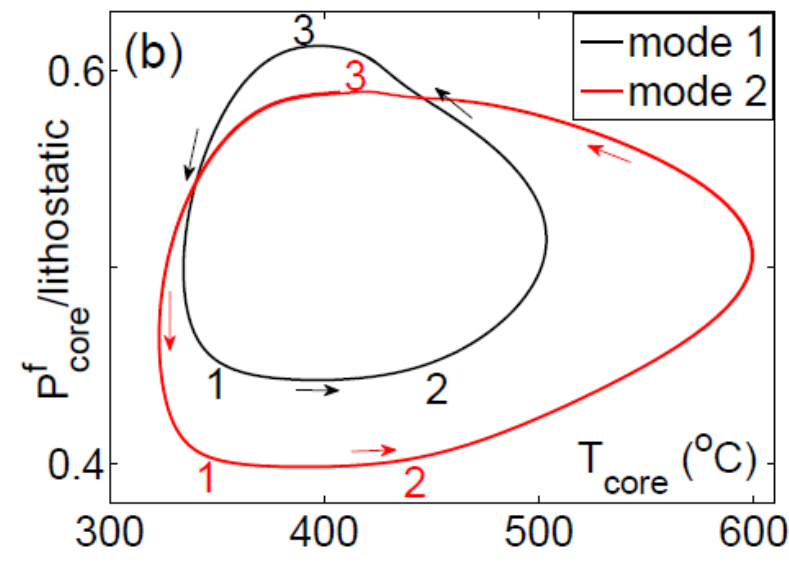
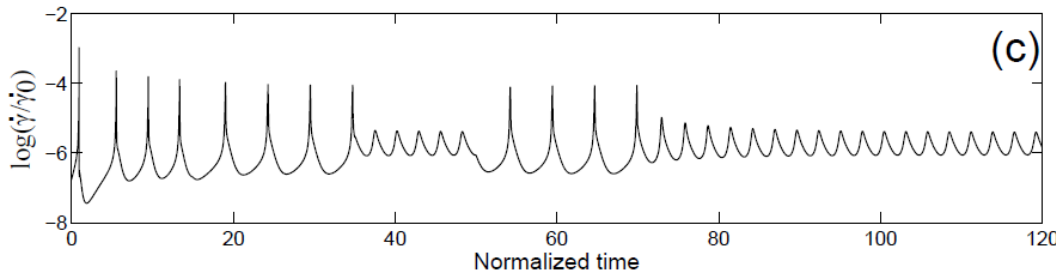
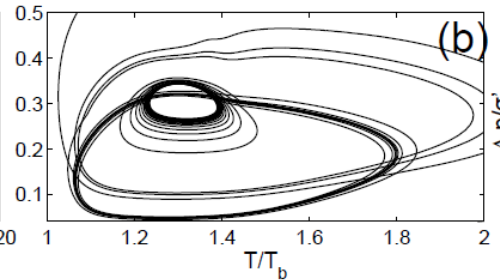
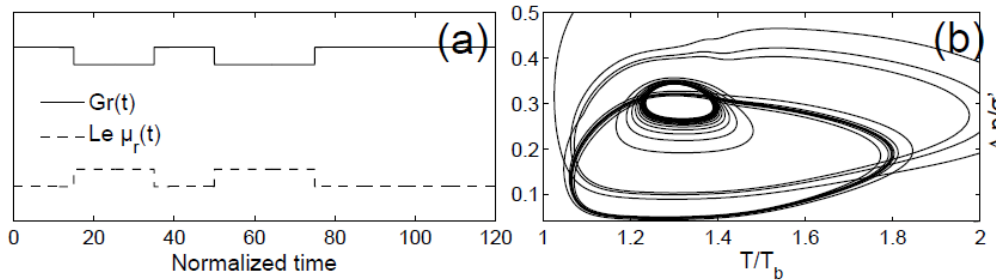
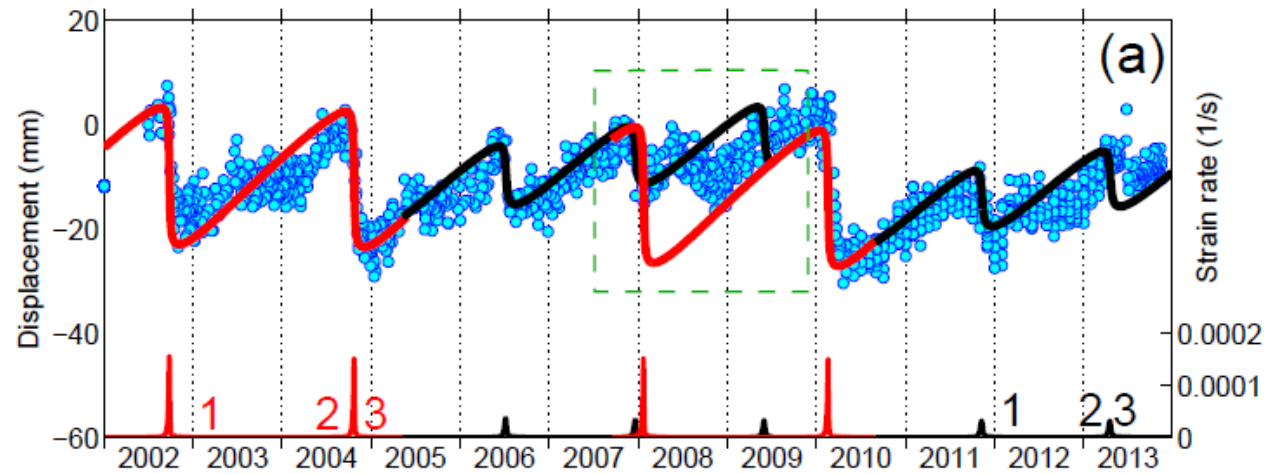
Cascadia earthquake sources



Rogers & Dragert, Science (2003)

Alevizos et al, JGR (2014)

Chaotic signals – Gisborne (New Zealand)



Matching length scales

- Intern. Ocean Discovery Program
- Japan Trench Fast Drilling Project



Nature news, Dec 2013:

“The localization of deformation onto a limited thickness (~5 meters) of pelagic clay is the defining characteristic of the shallow earthquake fault” (Chester et al / Science 2013). “**That’s just weird**” says Emily Brodsky (UC Santa Cruz)

Table 3. Material Parameters Inverted From the ETS Sequences, After Fitting the GPS Data^a

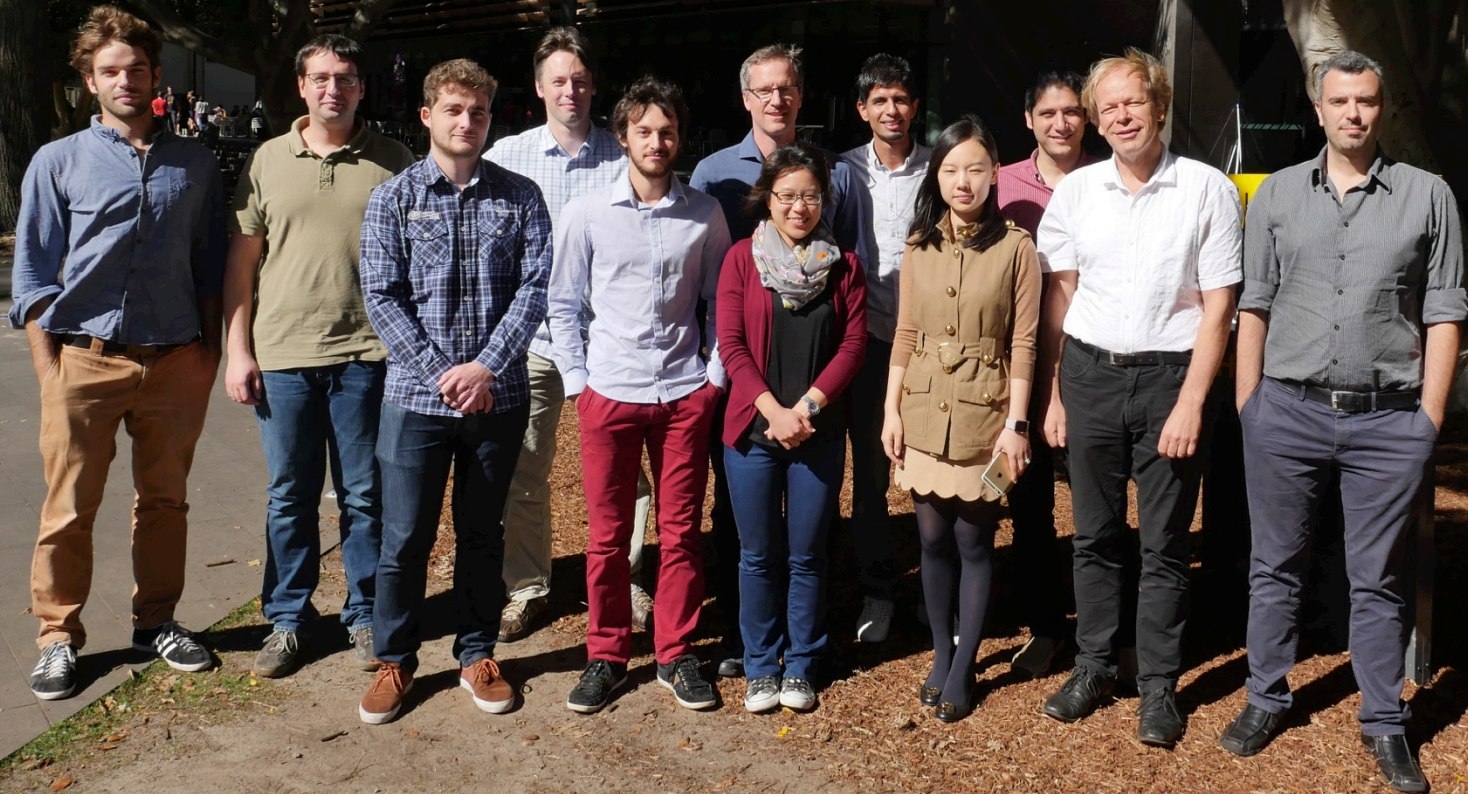
| Parameter | Units | ALBH | GISB 1 | GISB 2 |
|------------------------|----------|-----------|-------------|-------------|
| $\dot{\gamma}_0$ | s^{-1} | 200 | 230 | 230 |
| d | m | 6.4 | 6.4 | 6.4 |
| $\bar{\sigma}'_n$ | MPa | 49 | 49 | 74 |
| $\beta_T \bar{\tau}_n$ | MPa | 0.3 | 0.26 | 0.20 |
| k_F | s^{-1} | 10^8 | 10^8 | 10^8 |
| Q_F | kJ/mol | 114 | 114 | 114 |
| k_R | s^{-1} | 10^{-2} | 10^{-2} | 10^{-2} |
| ΔH | kJ/mol | 80 | 80 | 80 |
| Q_R | kJ/mol | 34 | 34 | 34 |



In Conclusion



Thank you!



<http://ugg.unsw.edu.au>

July 5-7, 2017 Paris, France

The 6th International Conference on
Coupled Thermo-Hydro-Mechanical-
Chemical (THMC) Processes in
Geosystems: Multiphysical instabilities
across the scales



UNSW



Important dates

Abstract submission: 15 November 2016

Abstract acceptance: 15 December 2016

Early-bird registration: 31 January 2017

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